

# Use of the Near Detector in the MINOS $\nu_e$ appearance measurement

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DUSEL Beamline WG Meeting  
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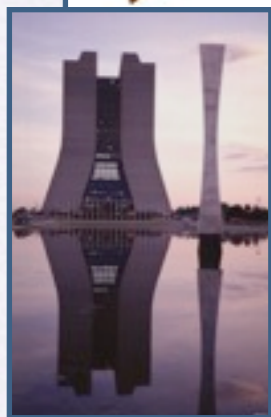
# Using the MINOS ND

- Talk with different focus: the analysis experience.
  - Many slides from last Friday's W&C.
- The main analyses in MINOS use the Near Detector for several important things:
  - Relative calibration.
  - Beam uncertainties.
  - Measurement of the signal before oscillations.
- The  $\nu_e$  appearance analysis is no different except it uses it to measure the background instead of the signal.

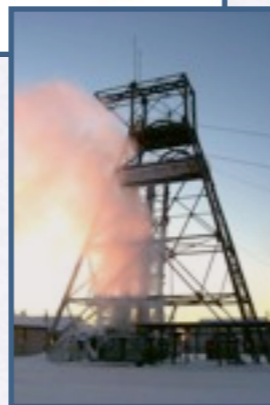
# MINOS in a nutshell

- Produce a high intensity beam of muon neutrinos at Fermilab.
- Measure background at the Near Detector and use it to predict the Far Detector spectrum.
- If neutrinos oscillate we will observe a distortion in the data at the Far Detector in Soudan, 735 km away.

## Main Injector Neutrino Oscillation Search

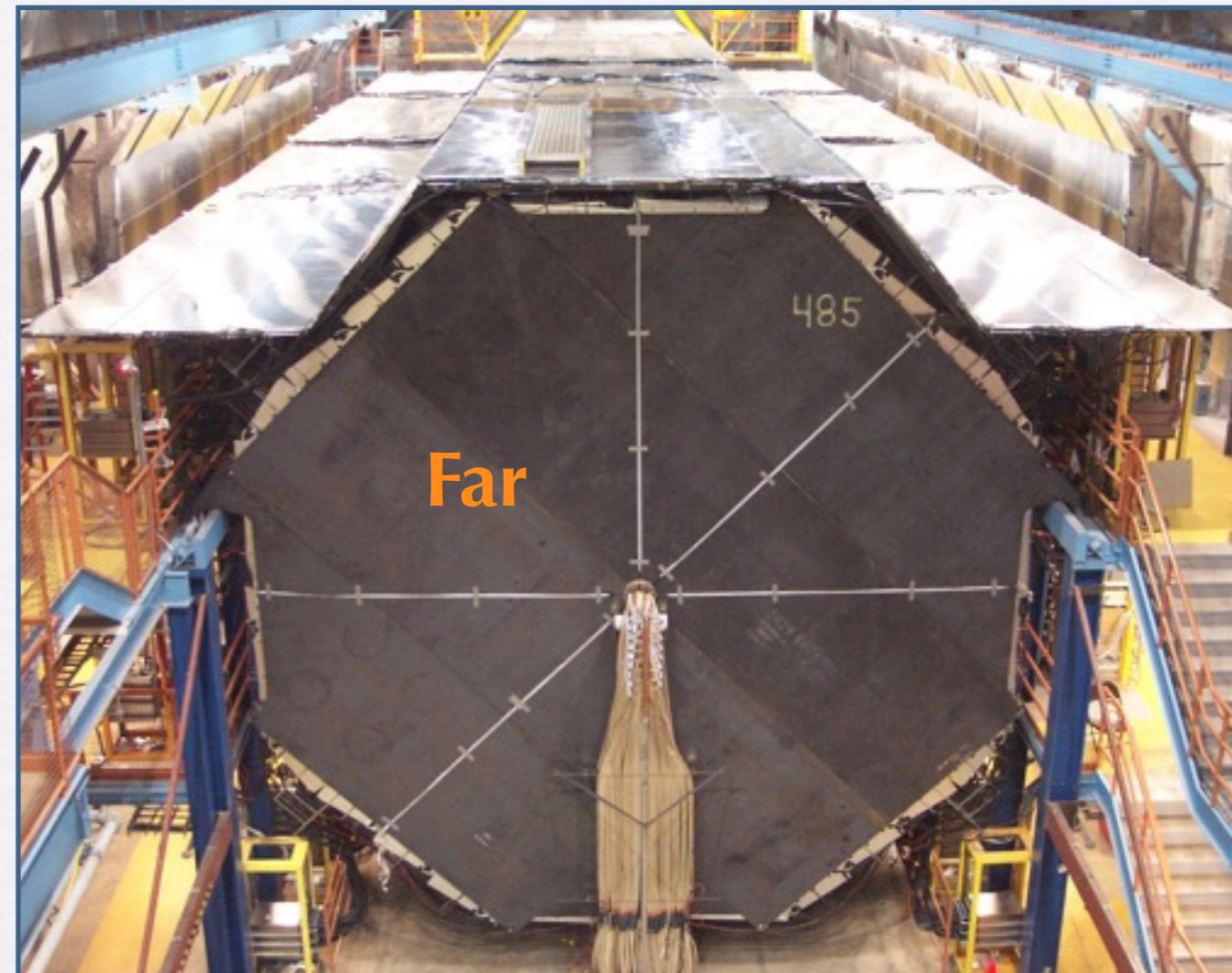


← long baseline →  
concept



# The MINOS detectors

- Functionally identical: **Near and Far detectors**
- 1 inch thick octagonal steel planes, alternating with planes of 4.1 cm x 1 cm scintillator strips, up to 8m long. Magnetized.
  - **Near:** ~ 1kton, 282 steel squashed octagons. Partially instrumented.
  - **Far:** 5.4 kton, 486 (8m/octagon) fully instrumented planes.



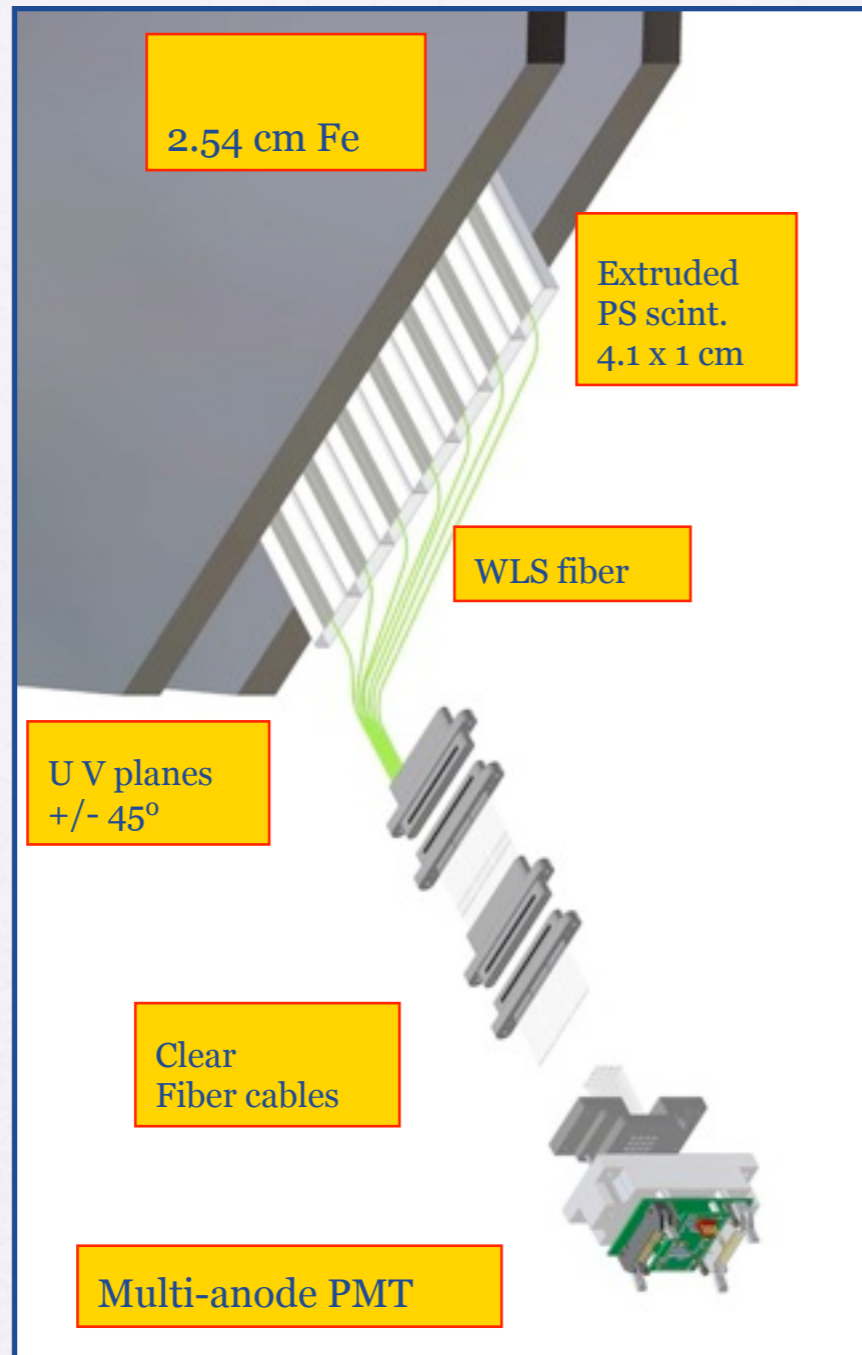
# MINOS detector technology



M16

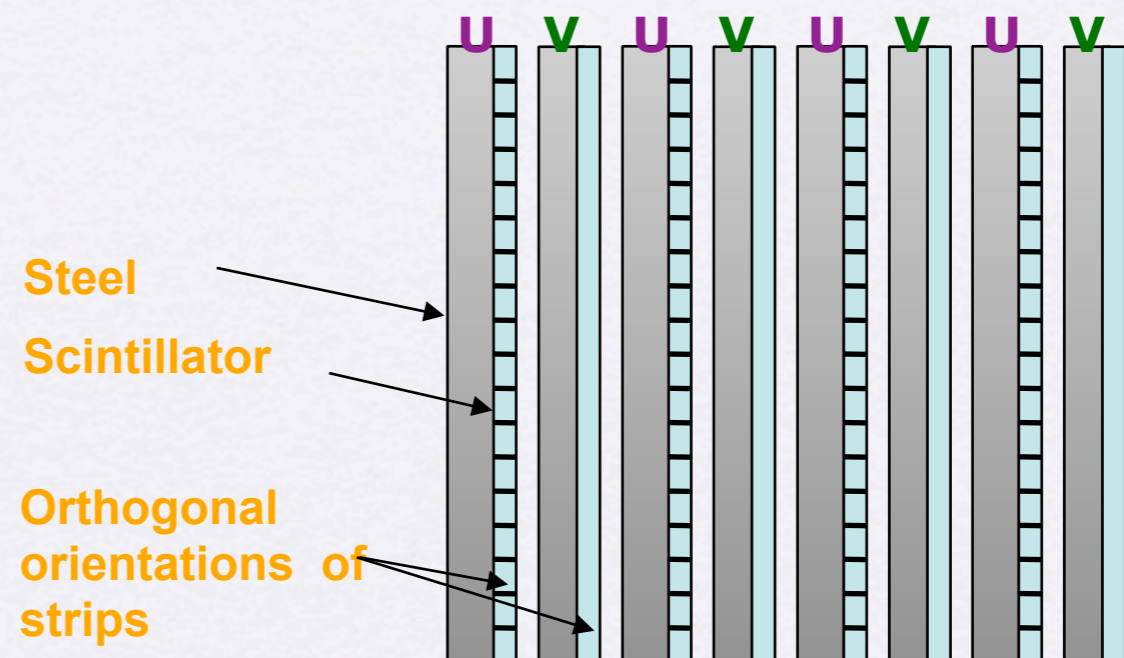


M64

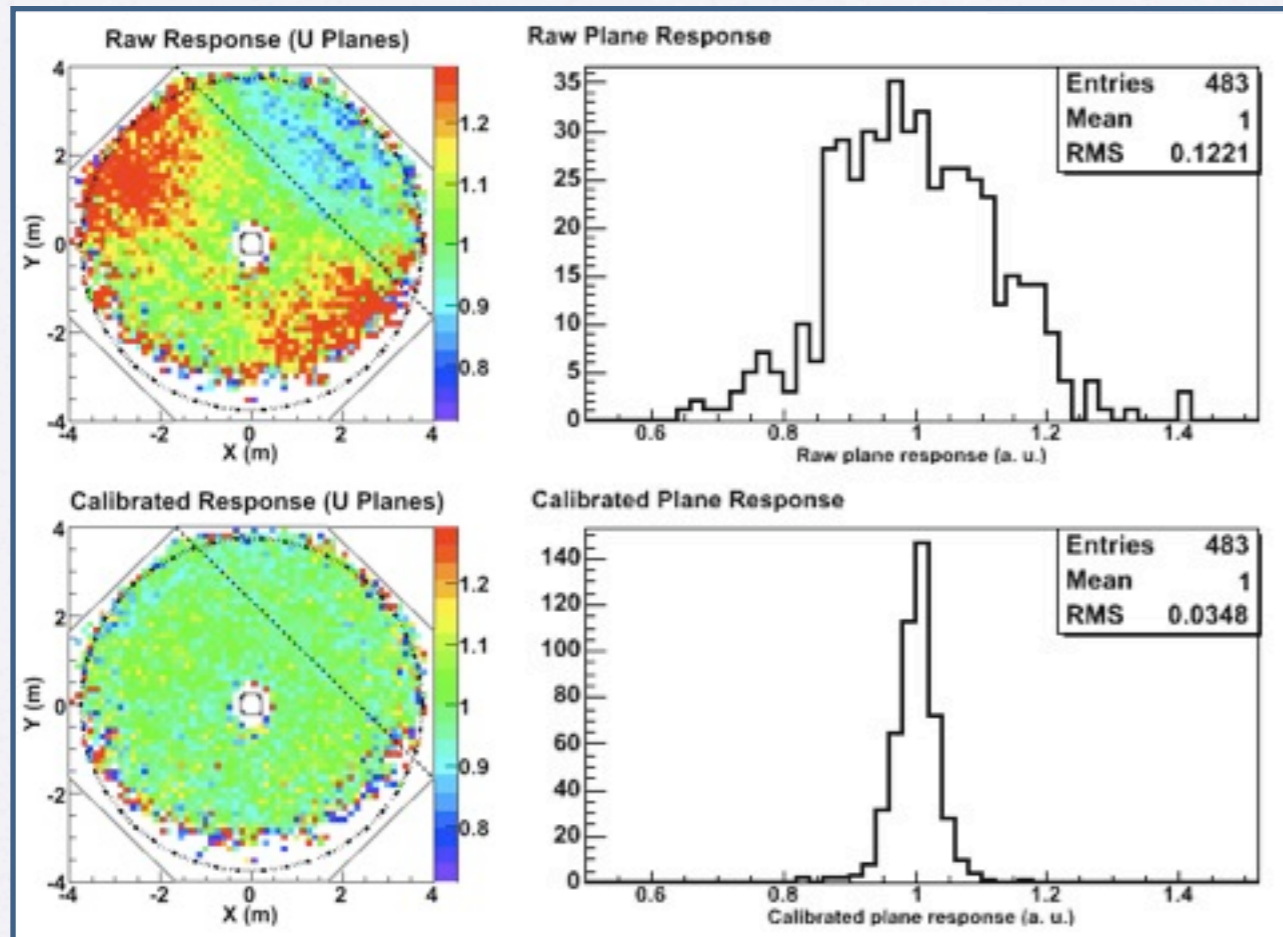


- **In both detectors:**

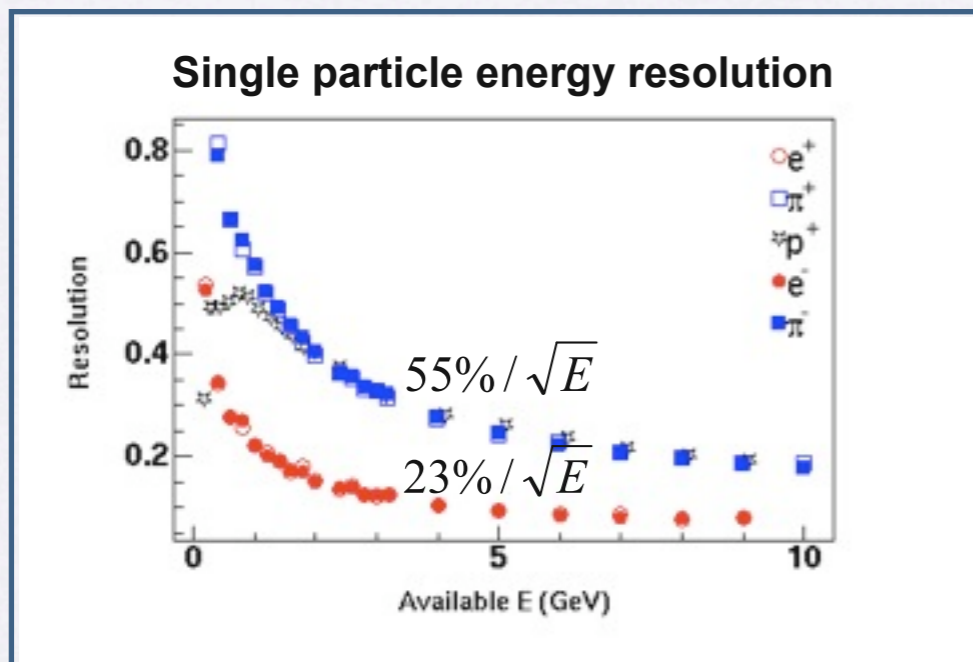
- co-extruded polystyrene scintillator strips
- orthogonal orientation on alternate planes - U, V
- optical fiber readout to multi-anode PMTs (M64 for the ND, M16 for the FD)



# MINOS calibration system

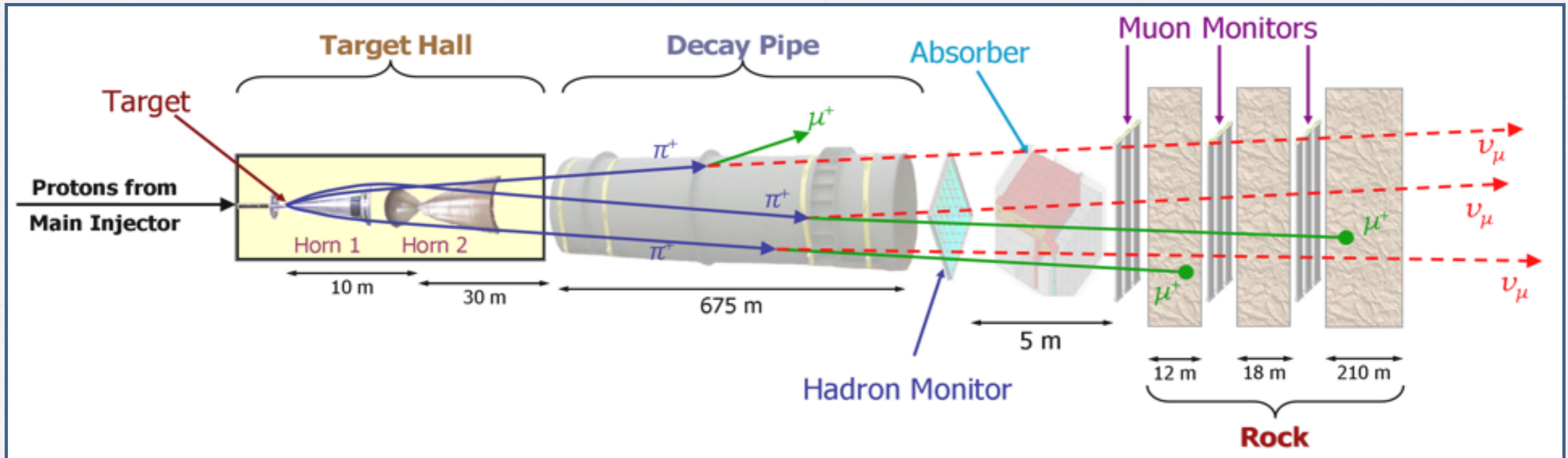


- Calibration of ND and FD response using:
  - Light Injection system (PMT gain)
  - Cosmic ray muons (strip to strip and detector-to-detector)
  - Calibration detector (overall energy scale)
    - mini-Minos in a CERN test beam (CalDet)

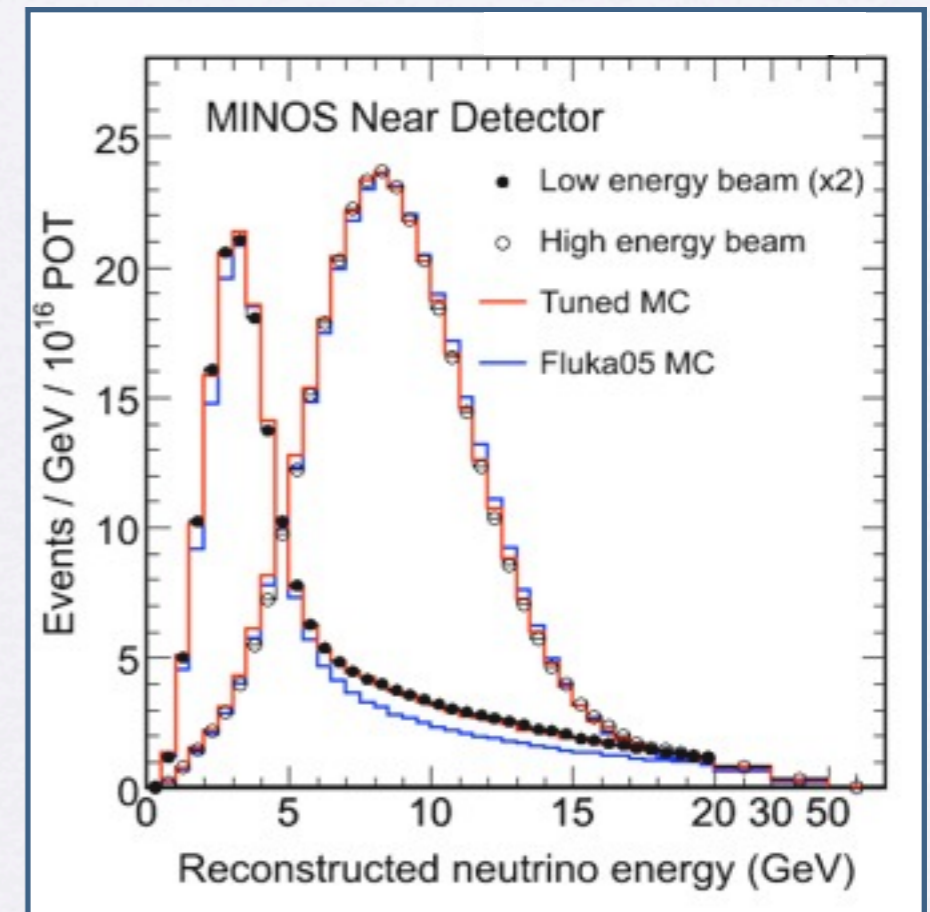


- Energy scale calibration:
  - 3.1% relative error in ND
  - 2.3% relative error in FD
  - ND-FD relative: 3.8%

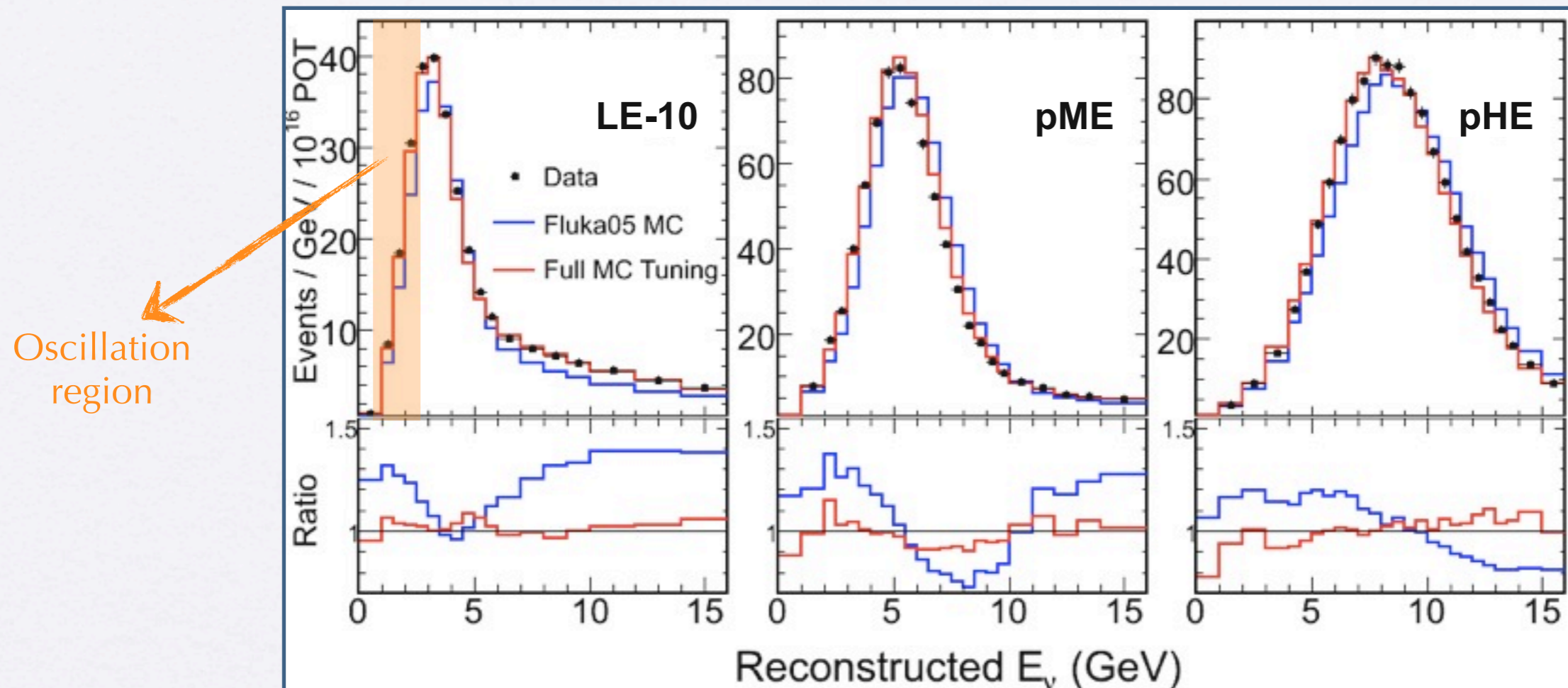
# Producing neutrinos



- Neutrinos from the Main Injector (NuMI)
- 10  $\mu\text{s}$  spill of 120 GeV protons every 2.2 s
- Currently 275 kW typical beam power
- Currently  $3.0 \times 10^{13}$  protons per pulse
- Neutrino spectrum changes with target position.



# Reconstructed Beam Spectrum



-10 cm, 185 kA

-100 cm, 200 kA

-250 cm, 200 kA

Discrepancies between data and Fluka05 Beam MC vary with beam setting: so source is due to beam modeling uncertainties rather than cross-section uncertainties.

MC tuned by fitting to hadronic  $x_F$  and  $p_T$  over 7 beam configurations (3 shown here).

- **Lesson: LE better understood, increase your flux below 1GeV.**

**LE-10 Data/tuned MC agree to ~5%.**

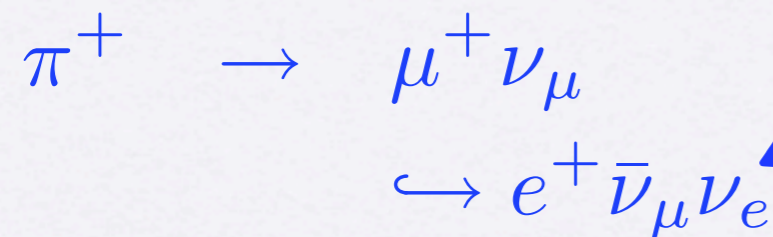
Worse agreement for higher energy beams.  
See Zarko's talk for details on this tuning and errors.



# Beam $\nu_e$ component

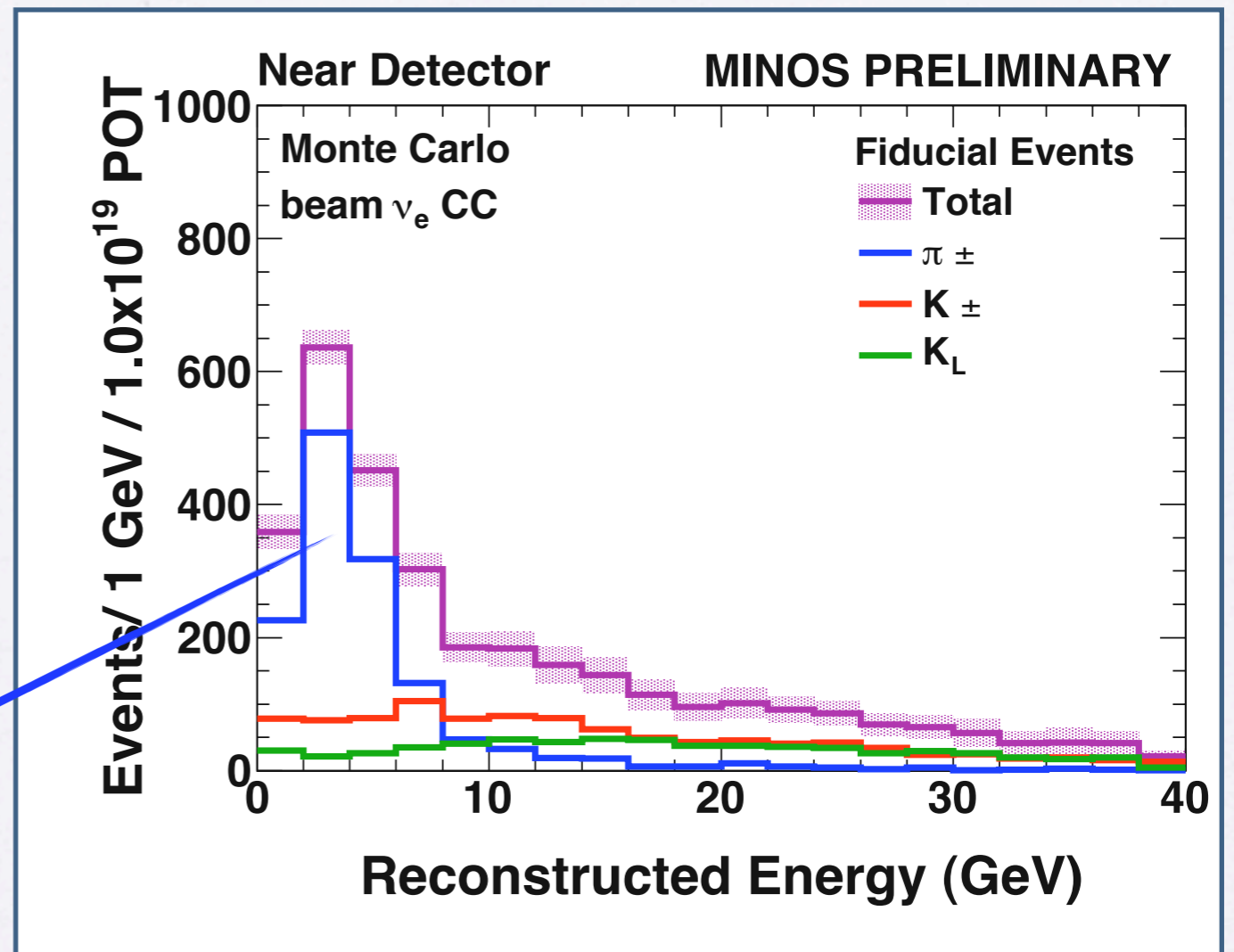
- Neutrino beam has 1.3% of  $\nu_e$  contamination from pion and kaon decays.

- Region of interest for the  $\nu_e$  oscillation analysis, 1-8GeV, dominated by events from secondary muon decays:



- Near and Far beam  $\nu_e$  spectra are constrained by using  $\nu_\mu$  events from several beam configurations.

- Uncertainties on the flux in the region of interest are  $\sim 10\%$ . After nue selection 9% in the Near, 13% in the Far.

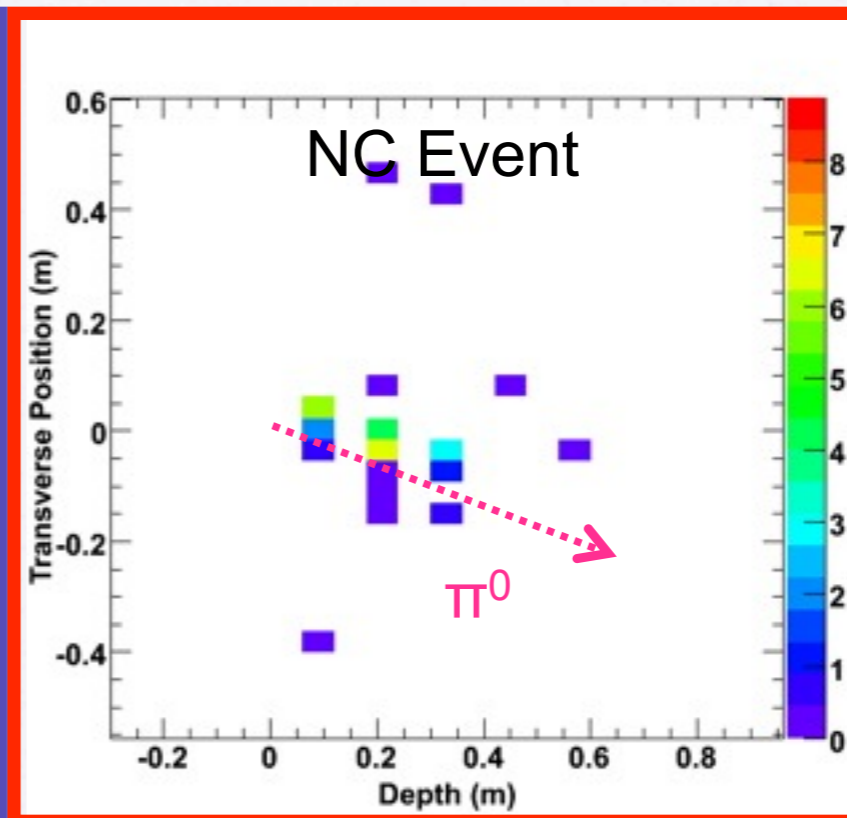
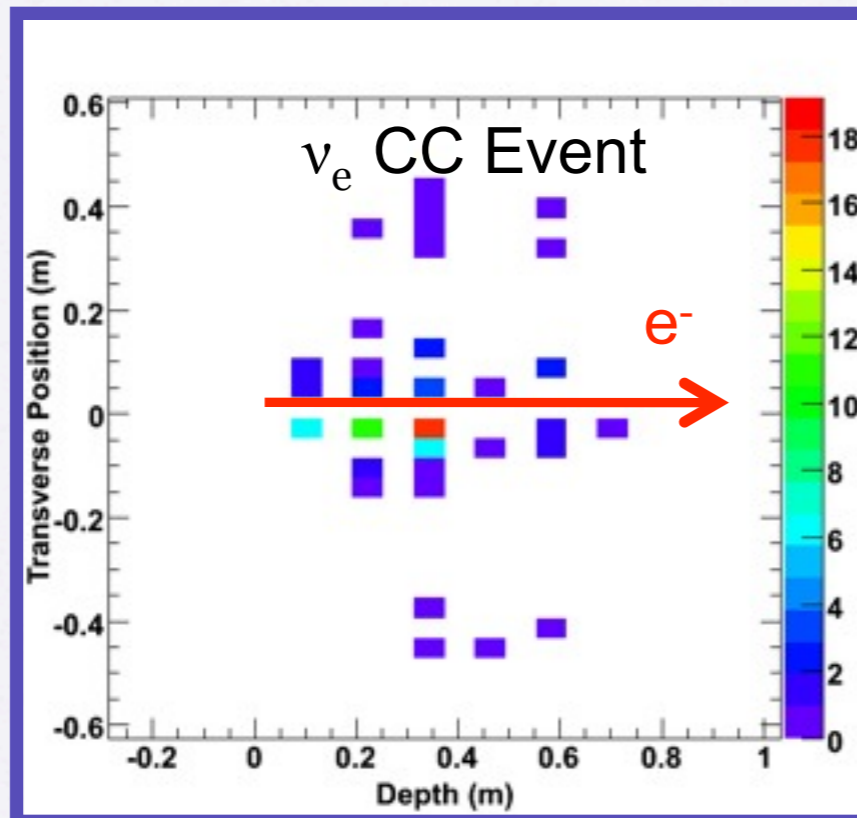


- Note off-axis beam would be dominated by kaons, less well known.
- Lessons:
- use same axis for the ND if you want to measure your beam  $\nu_e$ .
- if kaons are dominant you will need to constrain them separately, miniboone high?

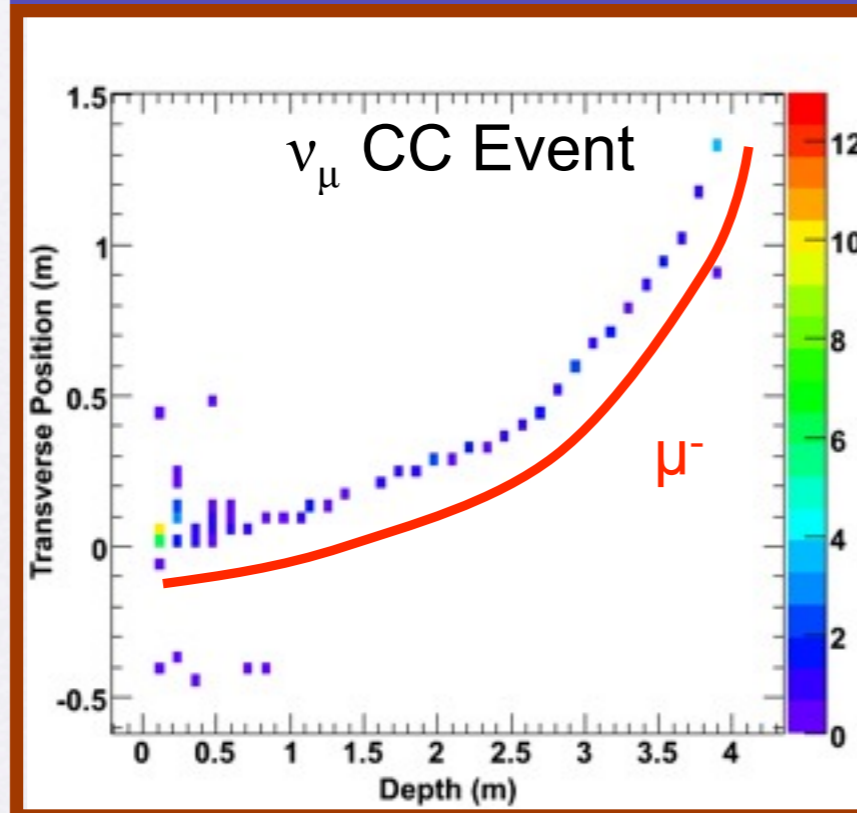
# Looking for $\nu_e$

To select  $\nu_e$  CC we focus on finding compact showers.

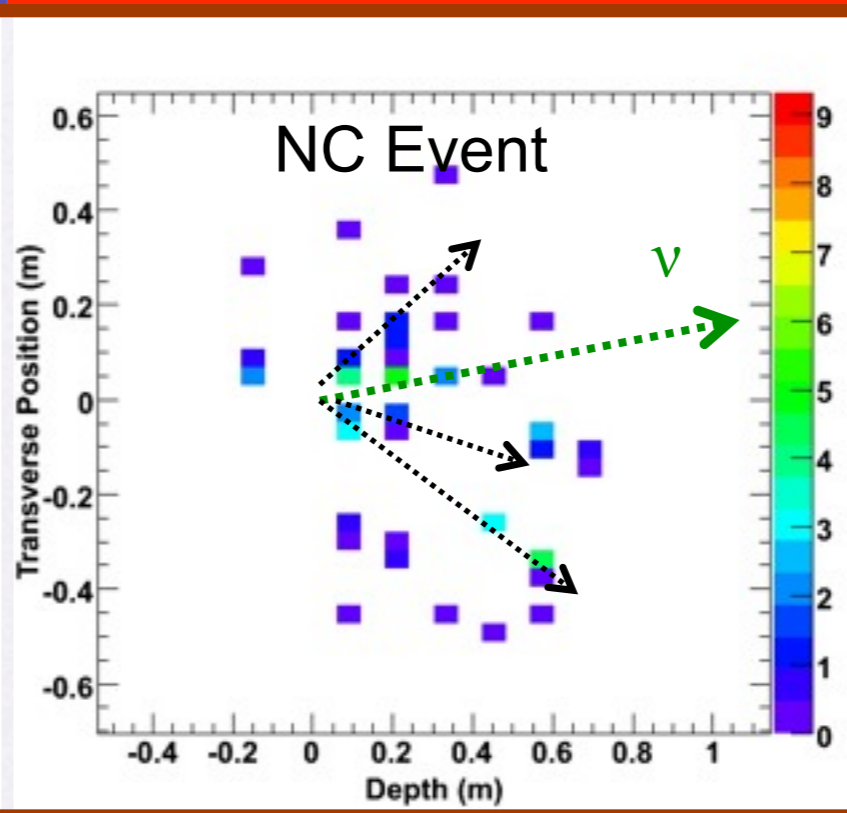
Signal



Irreducible Background

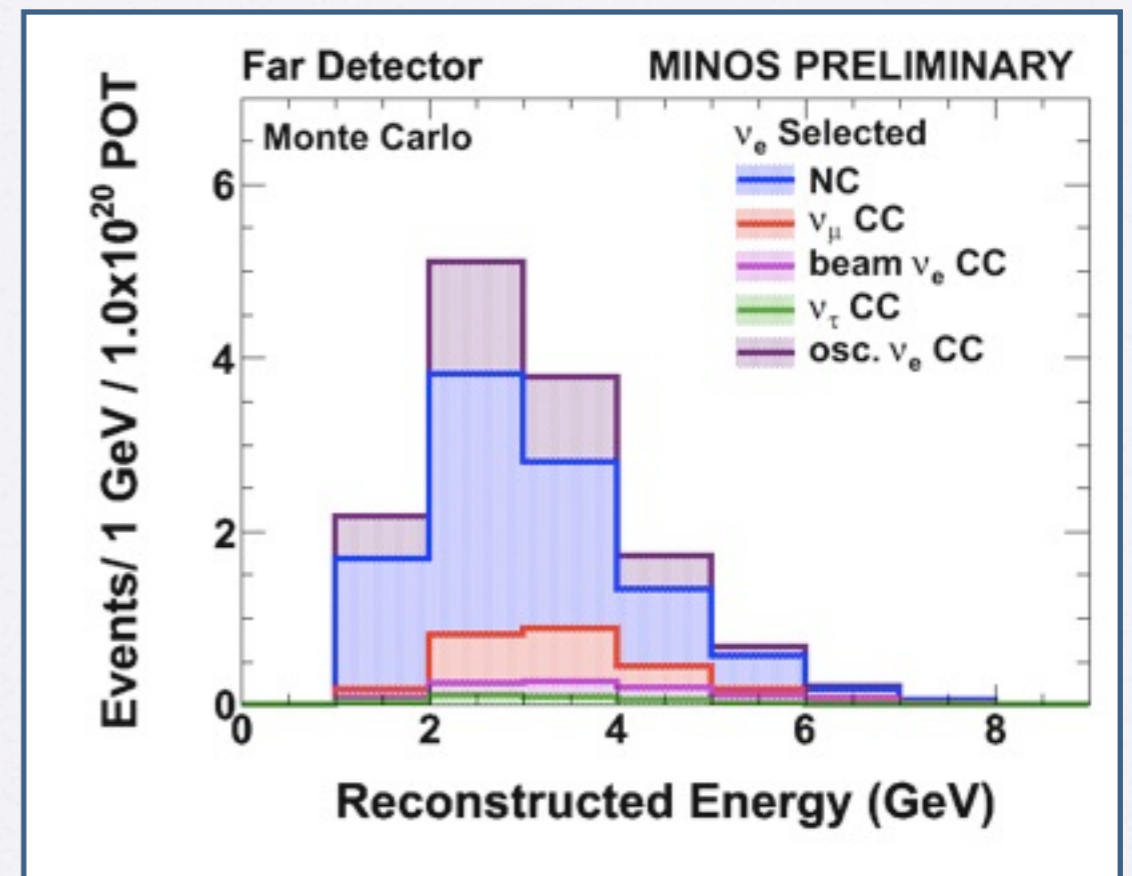
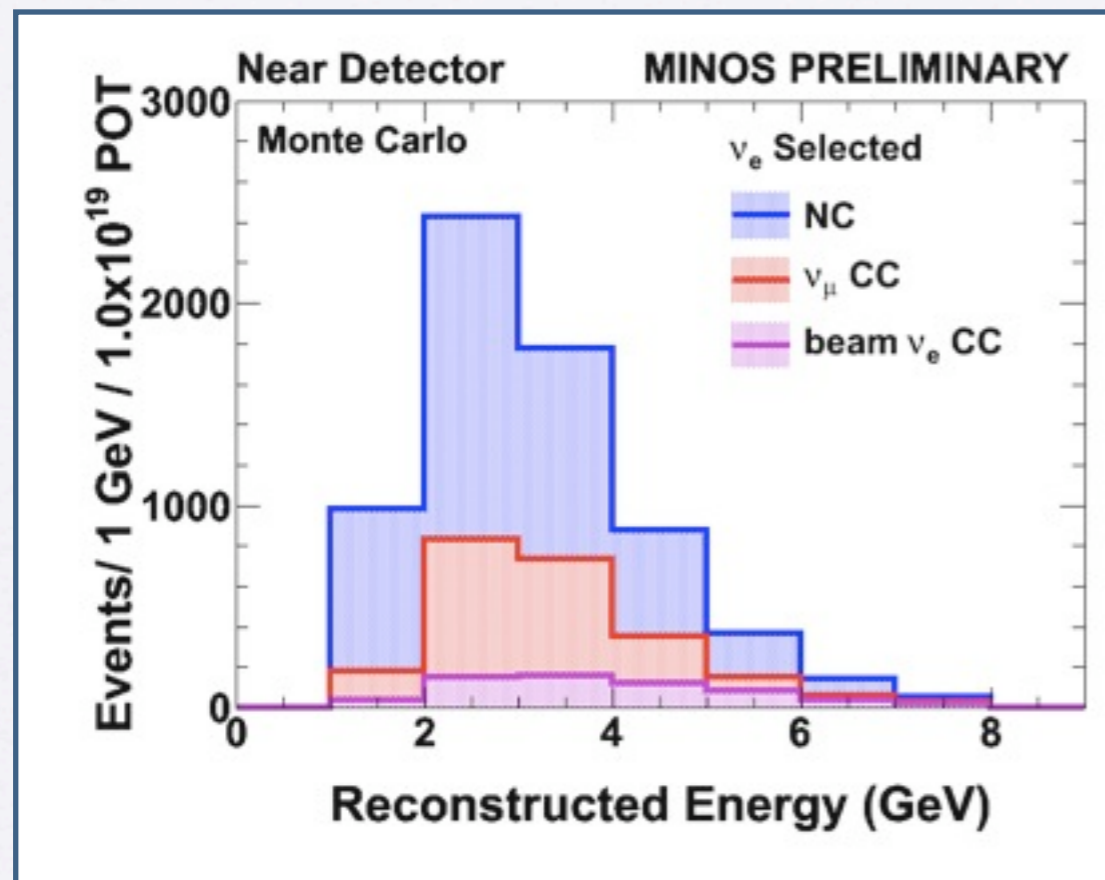


Reducible Background



Monte Carlo events

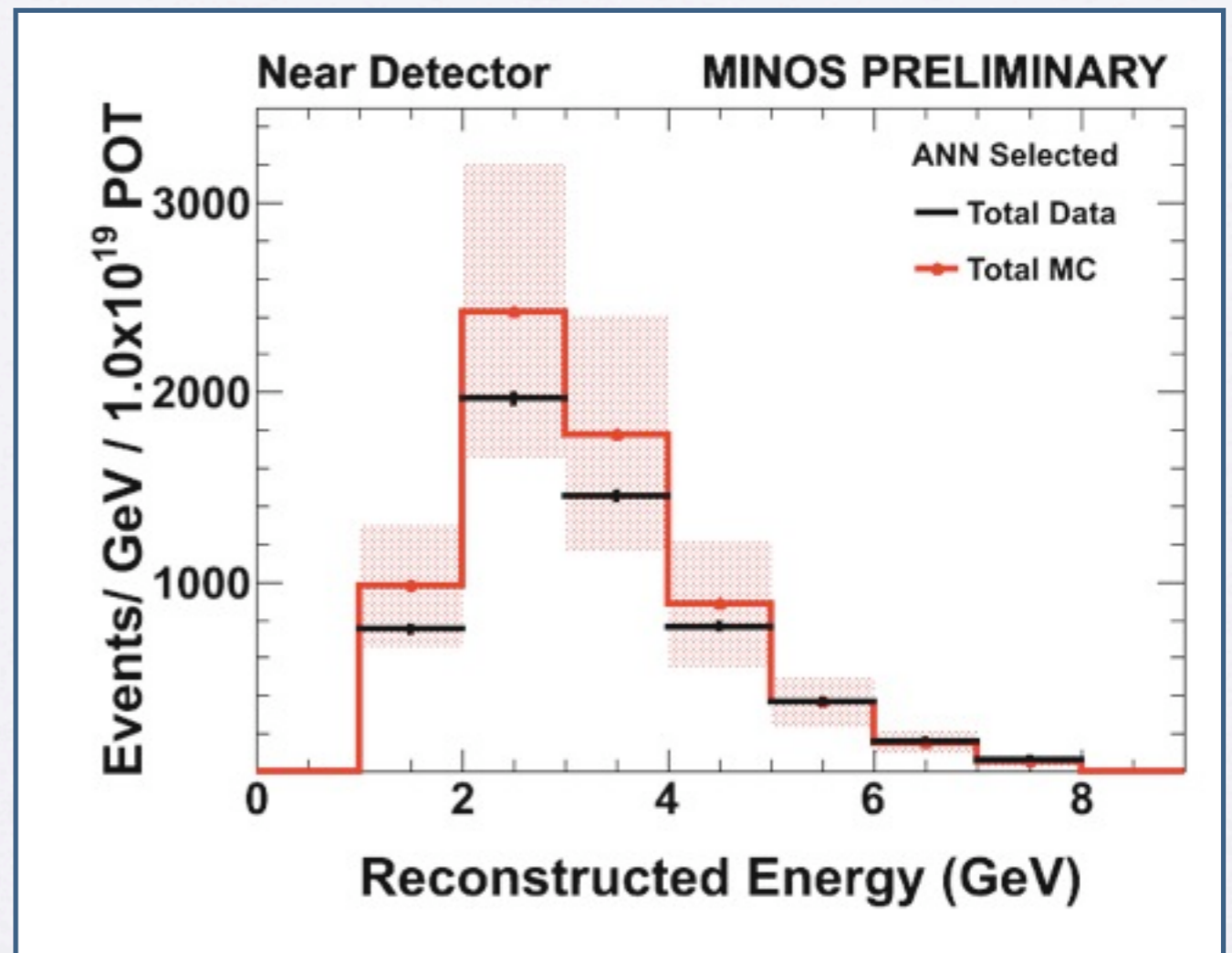
# $\nu_e$ appearance in MINOS



- When **selecting  $\nu_e$  event candidates** in the Near Detector we will have a mix of components that do not extrapolate in the same way to the Far Detector.
- We need to **separate the main background components** NC,  $\nu_\mu$  CC and beam  $\nu_e$  CC events, in the Near Detector.
- Then **extrapolate the background in the Far Detector** by extrapolating the components, oscillating the  $\nu_\mu$  CC component and calculating the  $\nu_\tau$  CC.
- Then **look for the  $\nu_e$  excess** arising from  $\nu_\mu$  to  $\nu_e$  oscillations in the Far Detector.
  - **Lesson: use same target mass, minimize Far/Near differences.**

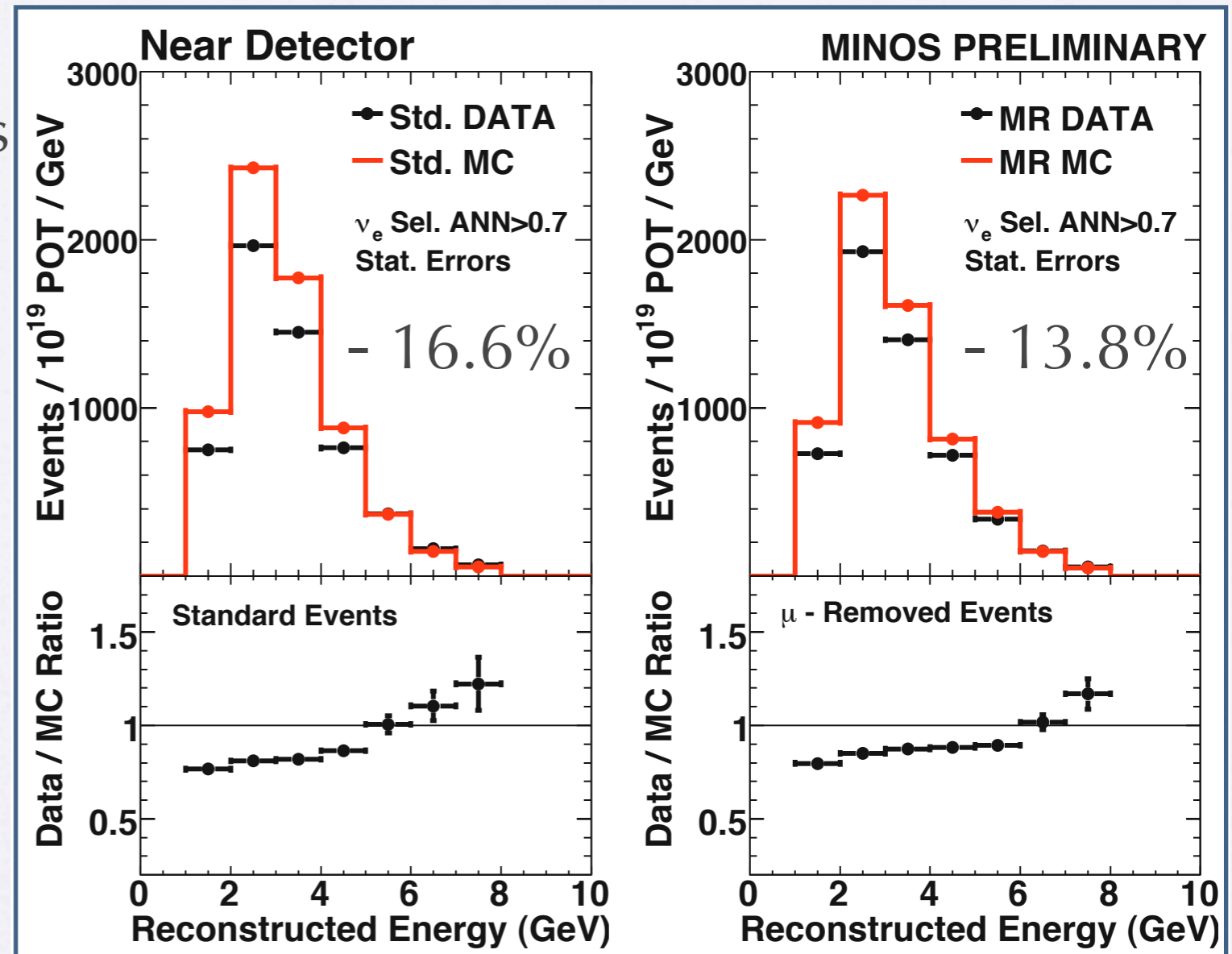
# $\nu_e$ selected Near Detector data

- MC tuned to external bubble chamber data for hadronization models.
  - External data sparse in our kinematic range.
  - Strong background rejection leaves just tails of distributions.
- It is not surprising that the data/MC shows disagreement with the model.
- Discrepancy is within the large uncertainties of the model.
- We have developed **two data-driven methods** to correct the model to match the data.
  - The MRCC method uses muon removed  $\nu_\mu$  CC to study the hadronic showers and correct MC.
  - The Horn on/off method uses the difference in background composition of the two horn configurations.
- **Lesson: Measure your background with same target mass, don't trust the MC.**



# Hadronic shower modeling in the $\nu_e$ selected data and muon-removed data

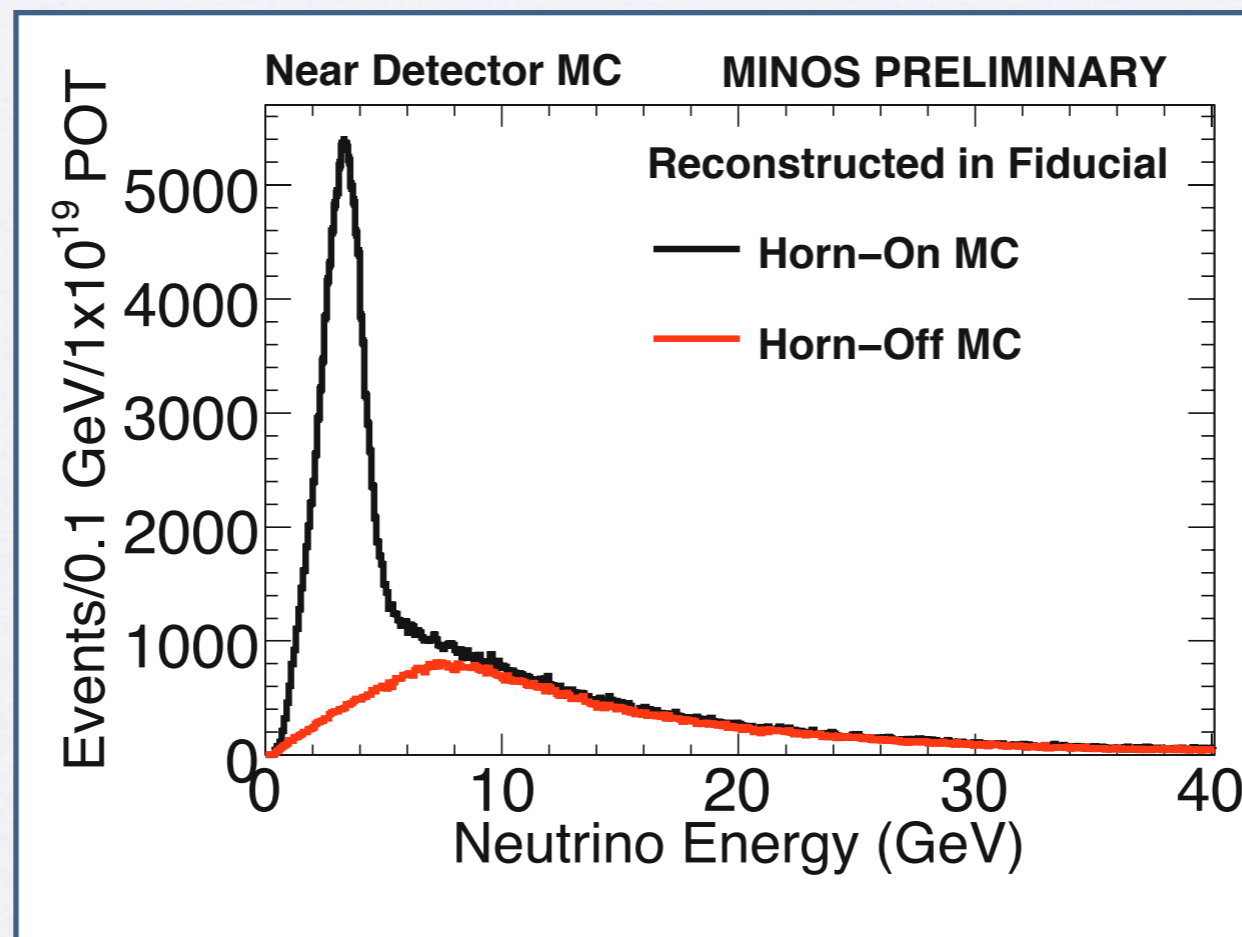
- We apply the  $\nu_e$  selection to the standard data and MC as well as to the Muon Removed data and MC.
- Discrepancy with the model shows the same trend not only in energy but in shower topology for both sets.
- Thus modeling of the hadronic shower is a major contribution to the disagreement.
- As the MRCC sample is independent, we can use it to obtain a **data-driven correction** to the model.



**Lesson: use the data creatively.**  
Note a relevant technique for a Water Cerenkov Detector or for a high resolution detector.

# Estimating the background using horn on and horn off data

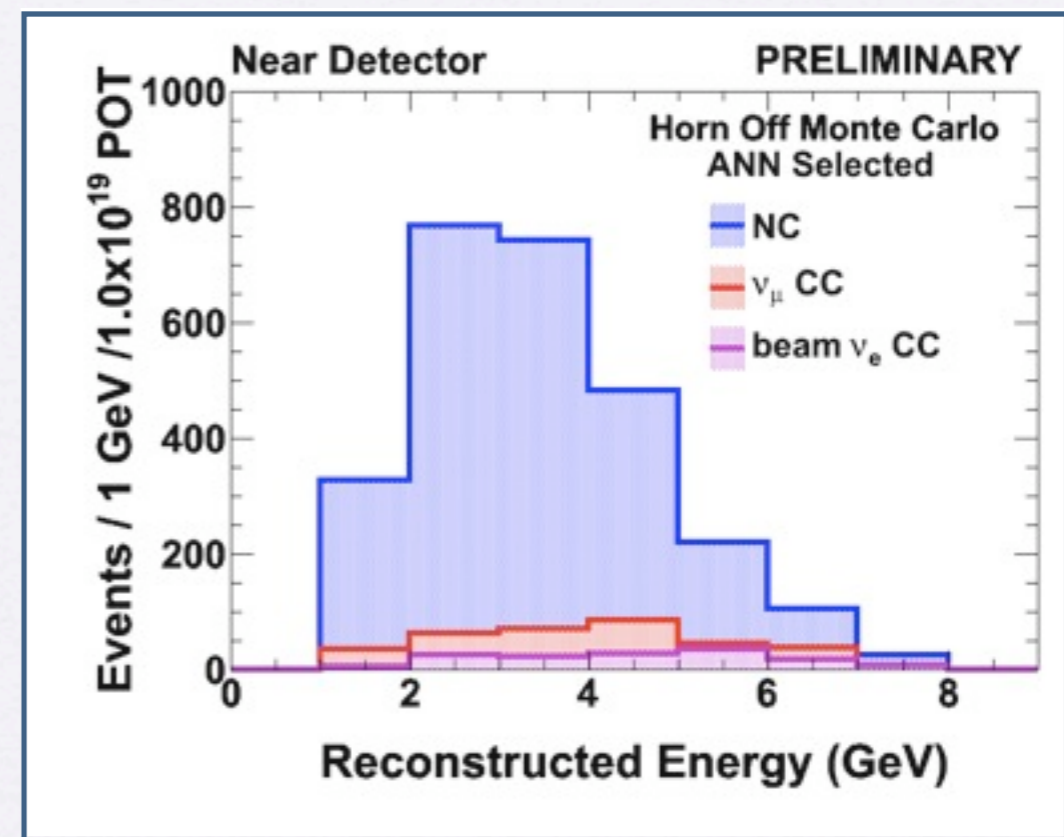
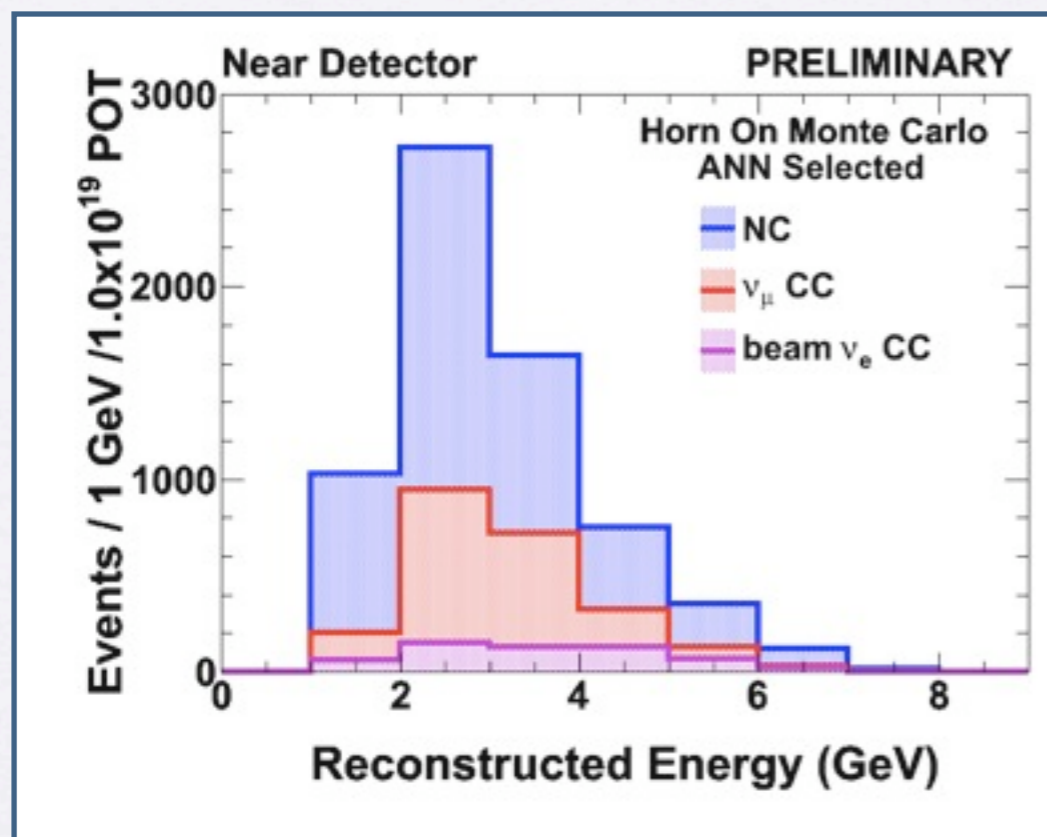
- When beam horns are turned off, the parent pions do not get focused, resulting in the disappearance of the low energy peak in the neutrino energy spectrum.



- The consequence is a spectrum dominated by NC arising from the long tail in true neutrino energy that gets measured in our region of interest in visible energy.

# Estimating the background using horn on and horn off data

- After applying the  $\nu_e$  selection cuts to the ND data, the composition of the selected events is thus very different with the NuMI horns on or off.



- Using the horn off spectrum which is dominated by NC, we can measure that component with better precision than in the horn on beam.
- **Lesson: Might be specific to MINOS resolution that we can do this, ie not sensitive to angular distribution of the showers.**

# Estimating the background using horn on and horn off data

- The **beam  $\nu_e$  flux** is obtained from the  $\nu_\mu$  CC flux which is constrained by data in the different beam configurations.
- The two main background components can be estimated using the number of data events in the horn on and horn off configurations:  $N^{\text{on}}$  and  $N^{\text{off}}$ .

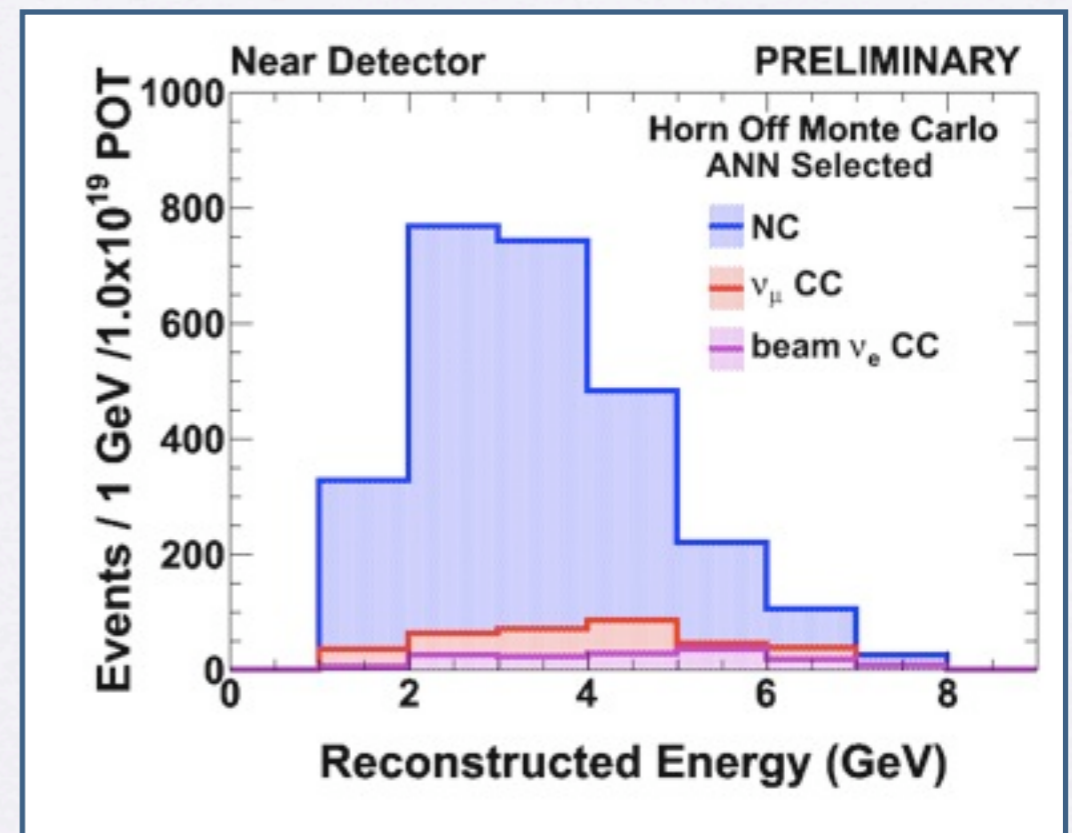
$$N^{\text{on}} = N_{\text{NC}} + N_{\text{CC}} + N_e \quad (1)$$

$$N^{\text{off}} = r_{\text{NC}} * N_{\text{NC}} + r_{\text{CC}} * N_{\text{CC}} + r_e * N_e \quad (2)$$

from MC:

$$r_{\text{NC(CC,e)}} = N_{\text{NC(CC,e)}}^{\text{off}} / N_{\text{NC(CC,e)}}^{\text{on}}$$

The key is to use the **Horn off/on ratios** for each component to solve:

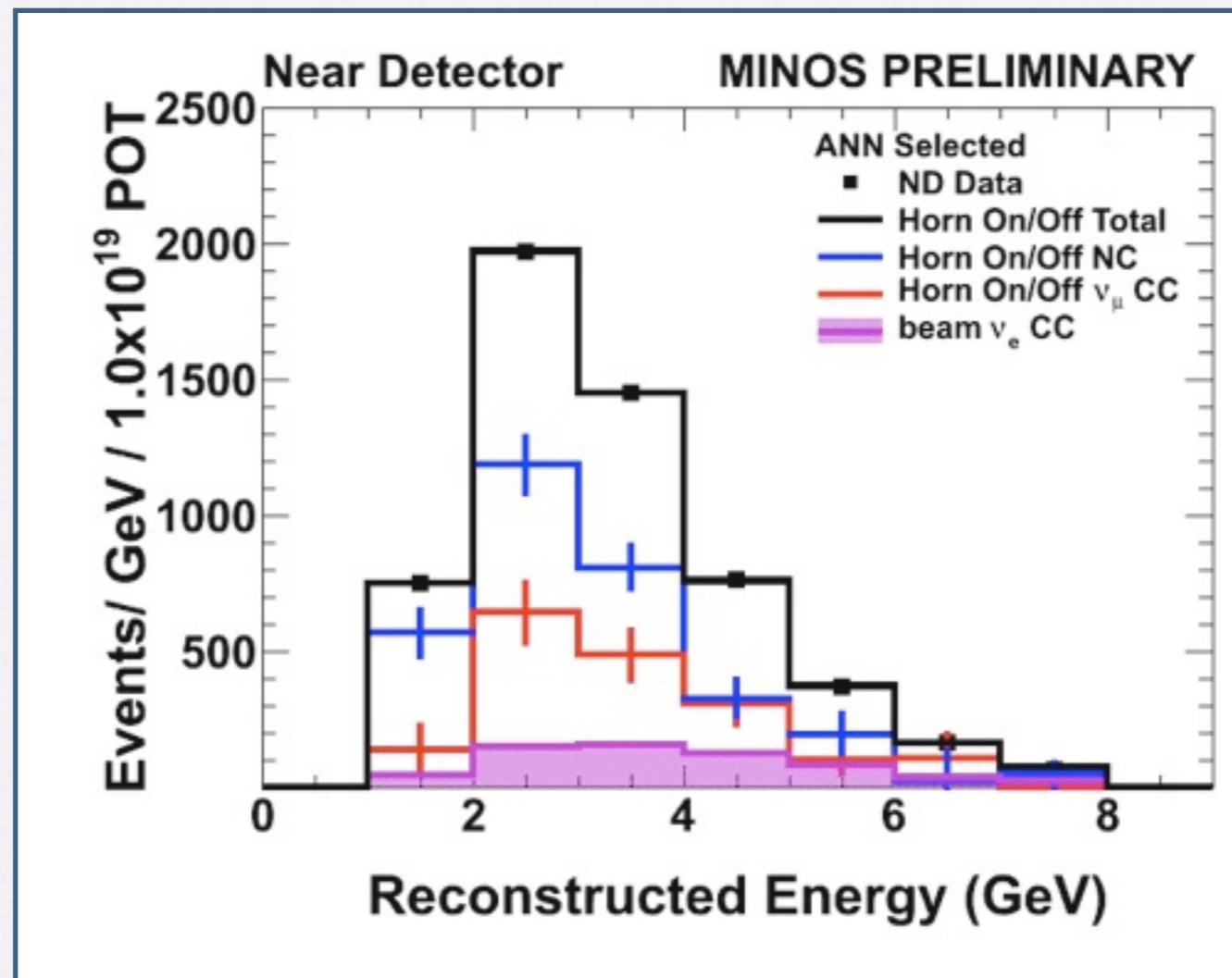


- Producing **data-driven predictions** for **NC** and  **$\nu_\mu$  CC** background for the horn on configuration.



# ND data-driven background

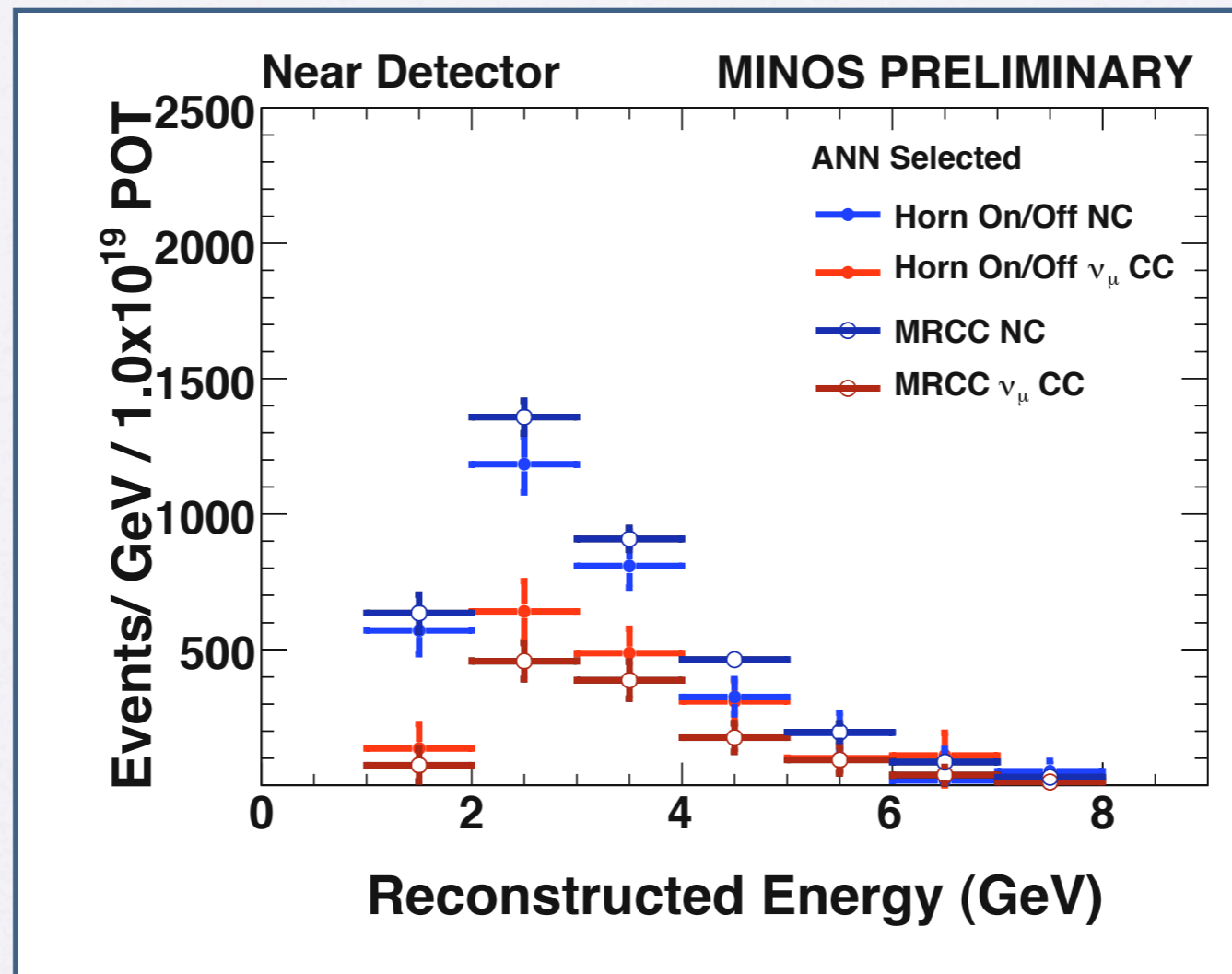
Results from the Horn on/off method



- The NC and  $\nu_\mu$  CC components for the the standard beam configuration are simultaneously solved in the horn on/off method and are by definition equal to the data after beam  $\nu_e$  subtraction.

# ND data-driven background

Results from the both methods



- The **two data-driven methods**, Horn on/off and MRCC, are in good agreement in the Near Detector NC and  $\nu_\mu$  CC background for the  $\nu_e$  analysis.
- Each background is then extrapolated to the Far Detector.

# ND data-driven background

Integral number of events selected

	Total	NC	$\nu_{\mu}$ CC	$\nu_e$ beam
MC	6764	4429	1742	593
Horn on/off	5524 $\pm$ 35	3150 $^{+292}_{-273}$	1781 $^{+366}_{-302}$	593 $\pm$ 178
MRCC		3674 $\pm$ 190	1236 $\pm$ 274	614 $\pm$ 186

scaled to  $1.0 \times 10^{19}$  POT

- The **two data-driven methods**, Horn on/off and MRCC, are in good agreement in the Near Detector NC and  $\nu_{\mu}$  CC background for the  $\nu_e$  analysis.
- Each background is then extrapolated to the Far Detector.

# ND data-driven background

## Horn off beam corrections

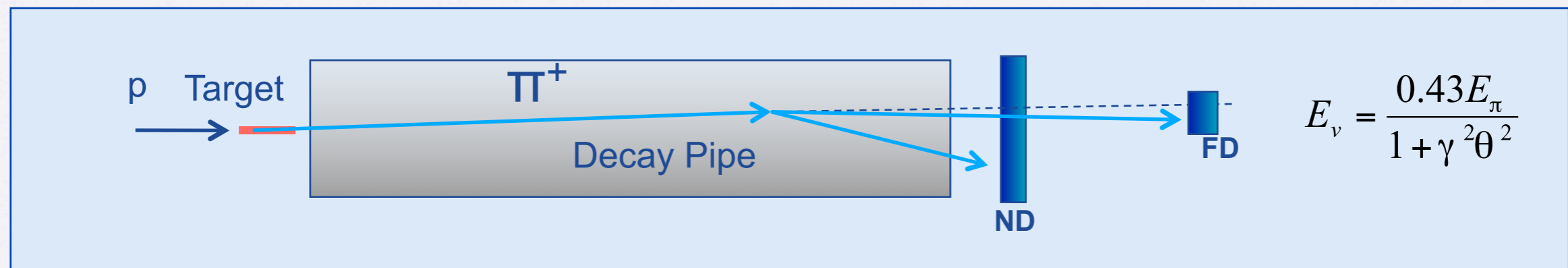
	Total	NC	$\nu_{\mu}$ CC	$\nu_e$ beam
MC	2680	2338	205	137
<b>Horn off</b>	$2105 \pm 62$	$1691^{+199}_{-182}$	$276^{+216}_{-148}$	$137 \pm 42$

scaled to  $1.0 \times 10^{19}$  POT

- $\text{NC (horn off)}/\text{NC (horn on)} = 0.54 \pm 0.08$
- NC corrected by 28% for horn off, 29% for horn on.

# Predicting the FD background

- Use Near Detector data to predict Far Detector spectrum.
- We expect the Far Detector spectrum to be similar to  $1/R^2$  scaled Near Detector spectrum, but not identical.



- Predict the event rate at each energy bin by correcting the expected Monte Carlo rate using the ratio of data to Monte Carlo in the Near Detector:

$$FD_i^{predicted} = \frac{FD_i^{MC}}{ND_i^{MC}} ND_i^{Data}$$

- The Monte Carlo provides necessary corrections due to energy smearing and acceptance.

# Other Far/Near differences

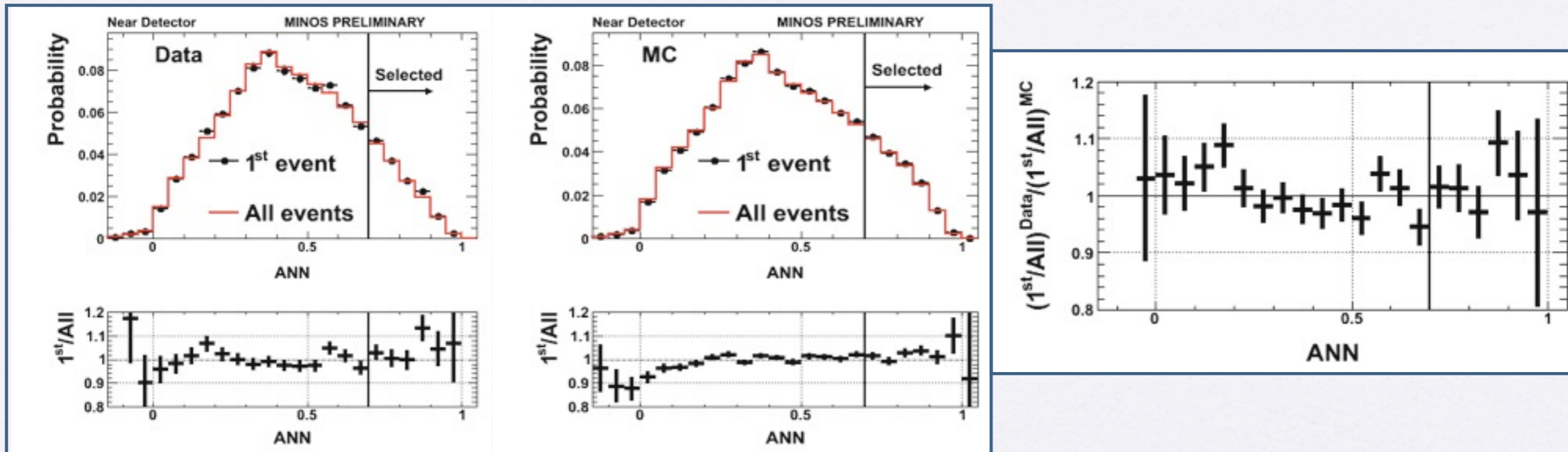
MINOS detectors are very similar, however there are small differences:

- Far/Near spectrum different due to **beamline geometry and oscillations** in the Far.
- **Readout patterns:**
  - Light level differences due to differences in fiber length.
  - Multiplexing in the Far (8 fibers per PMT pixel).
  - Partial (one-sided) readout in the Near.
- **Photomultipliers (M64 in Near Detector, M16 in Far):**
  - Different gains/front end electronics.
  - Different crosstalk patterns (also related to readout patterns).
- **Neutrino intensity:**
  - higher rates in the Near Detector thus faster readout.
- **Relative energy calibration.**

These considerations affect the Far/Near ratio and result in systematic errors.

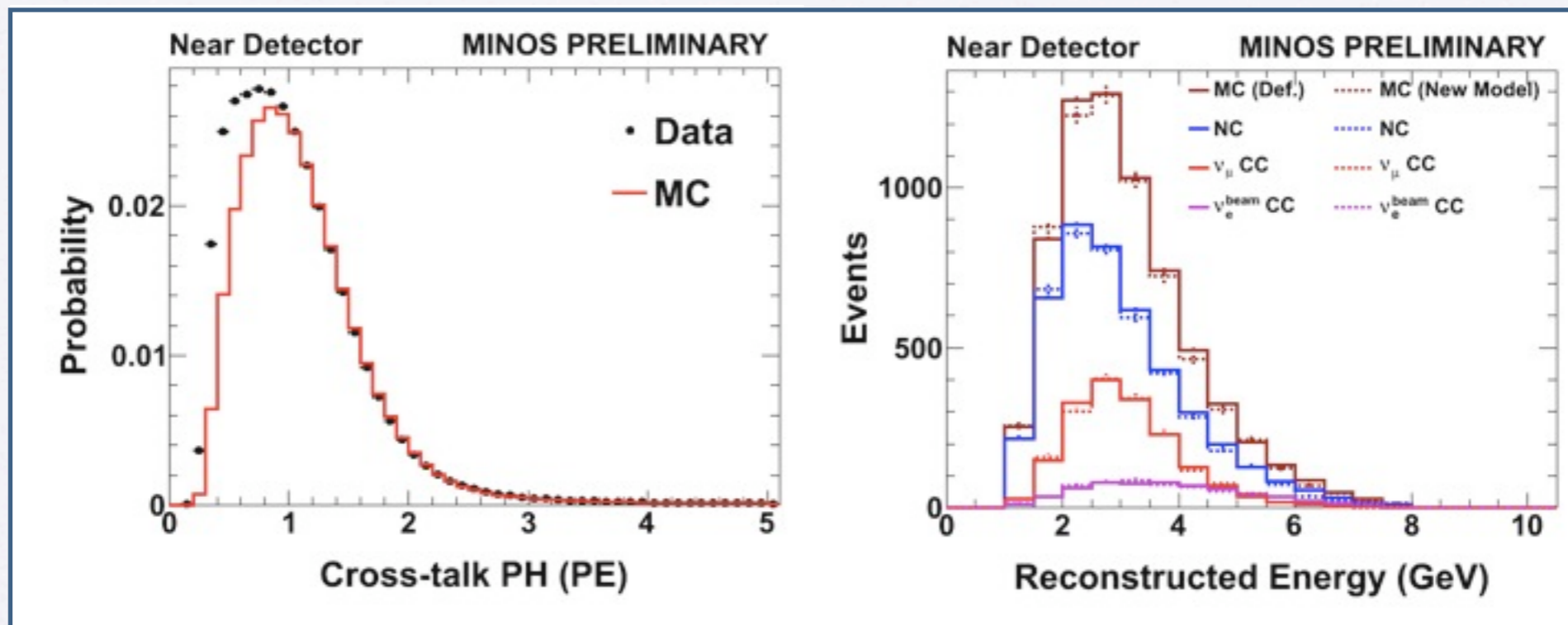
**Lesson: Make your detectors as similar as possible.**

# Intensity systematic



- Different rates at two detectors:  $\sim 8$  events in  $10\mu s$  spill window in ND and  $\sim 1$  event per day in FD.
- First event in each ND spill is unaffected by late activities of other events. We compare the 1st event to all other events to understand potential systematic effects.
- Difference in relative efficiencies between data and MC is taken as systematic error for the integral above the cut resulting in a 1% systematic error.
- **Lesson: Intensity will be more important for the water cerenkov detectors, move detector farther away? see 2km detector for T2k.**

# Crosstalk systematic

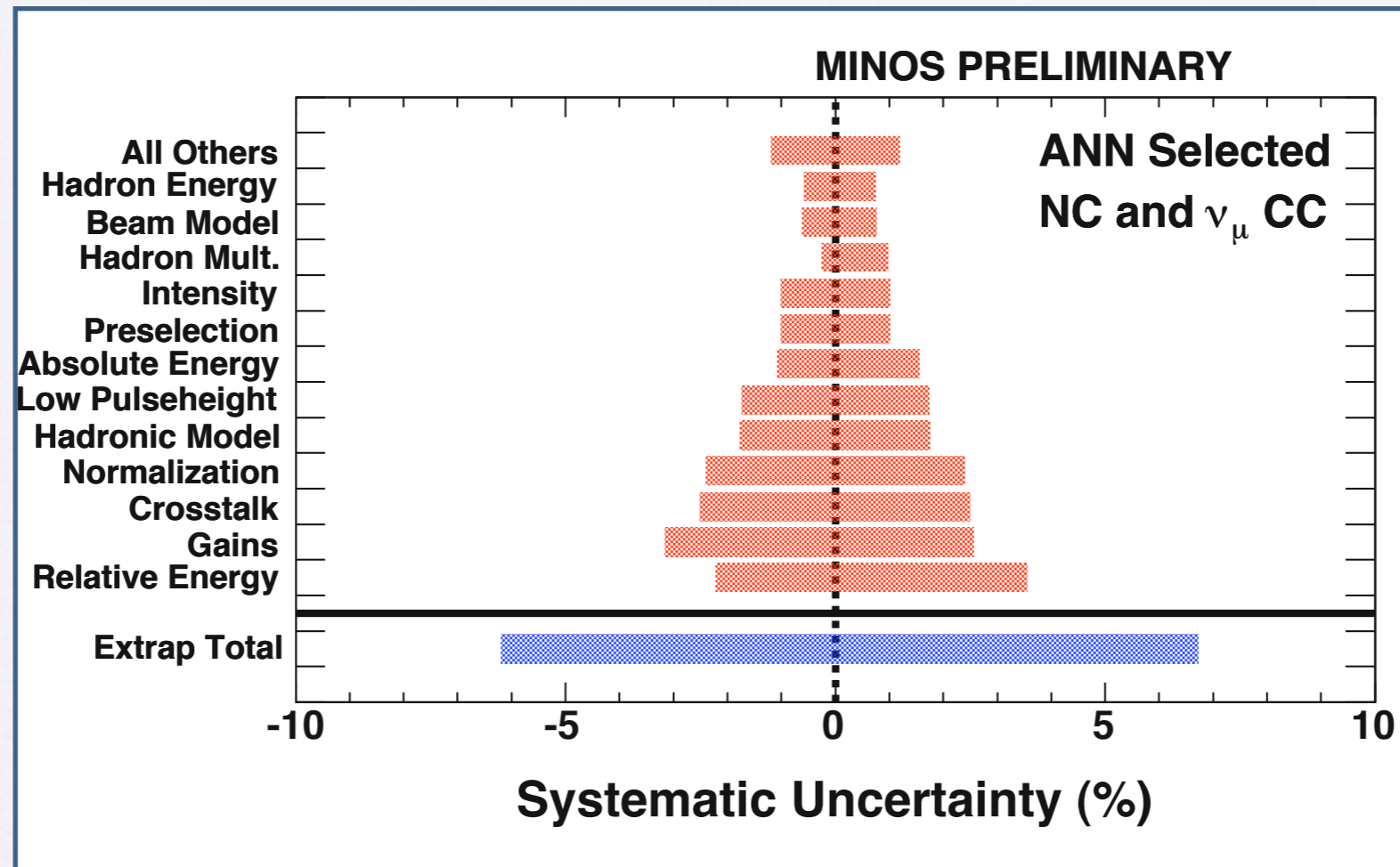


- PMT crosstalk is not well modeled in the Monte Carlo.
- Input variables to selections were constructed to avoid this problem by using hits greater than 2PE.
- The crosstalk model was improved by using cosmic ray muons.
- The difference between number of events selected in the current vs to the improved model was used as a systematic error.
- **Lesson: Use identical photodetectors, but you can't. Just as in MINOS there are limitations, need better pixelation.**



# FD background systematic errors

## Extrapolation errors



- For the main background components the larger systematics are **relative energy, gains, crosstalk and relative normalization**.
- **Lesson:** If everything is made equal, the only relevant parameters are the distance to both detectors, the mass and the efficiency at each of the detectors.

# FD background systematic errors

## Total errors

Preliminary Uncertainties	Horn On/Off
(1) Extrapolation	6.4%
(2) Systematic (separation method)	2.7%
(3) Statistical (separation method)	2.3%
Total (sum in quadrature)	7.3%
<b>Statistical error (data)</b>	<b>19%</b>

**Lesson: We can obtain a 7% systematic uncertainty because we have two similar detectors.**

# FD data-driven background

		Total	NC	$\nu_{\mu}$ CC	$\nu_{\tau}$ CC	$\nu_e$ beam
Data-driven Methods	Horn on/ off	27	18.2	5.1	1.1	2.2
	MRCC	28	21.1	3.6		

scaled to  $3.14 \times 10^{20}$  POT

- The two data-driven methods, Horn on/off and MRCC, are in excellent agreement in the Far Detector.
- $\sim 1$  event difference is well within errors.
- The horn on/off is the primary separation method.

**The background prediction at  $3.14 \times 10^{20}$  POT is:  
 $27 \pm 5(\text{stat}) \pm 2(\text{sys})$**

# $\nu_e$ appearance result:

MINOS PRELIMINARY

In case you missed it!

**Lesson: We want more statistics!**

$\nu_e$  appearance result:

**Observation 35 events**

**Expected Background  $27 \pm 5(\text{stat}) \pm 2(\text{sys})$**

for  $3.14 \times 10^{20}$  POT

MINOS PRELIMINARY

In case you missed it!

**Lesson: We want more statistics!**

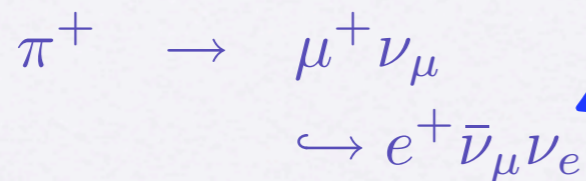
# Summary

- Having similar detectors provides invaluable measurement of the background when looking for electron neutrino appearance.
- Goal should be:
  - same target mass
  - similar intensity
  - similar photodetectors
- Any differences will be paid in systematic errors.

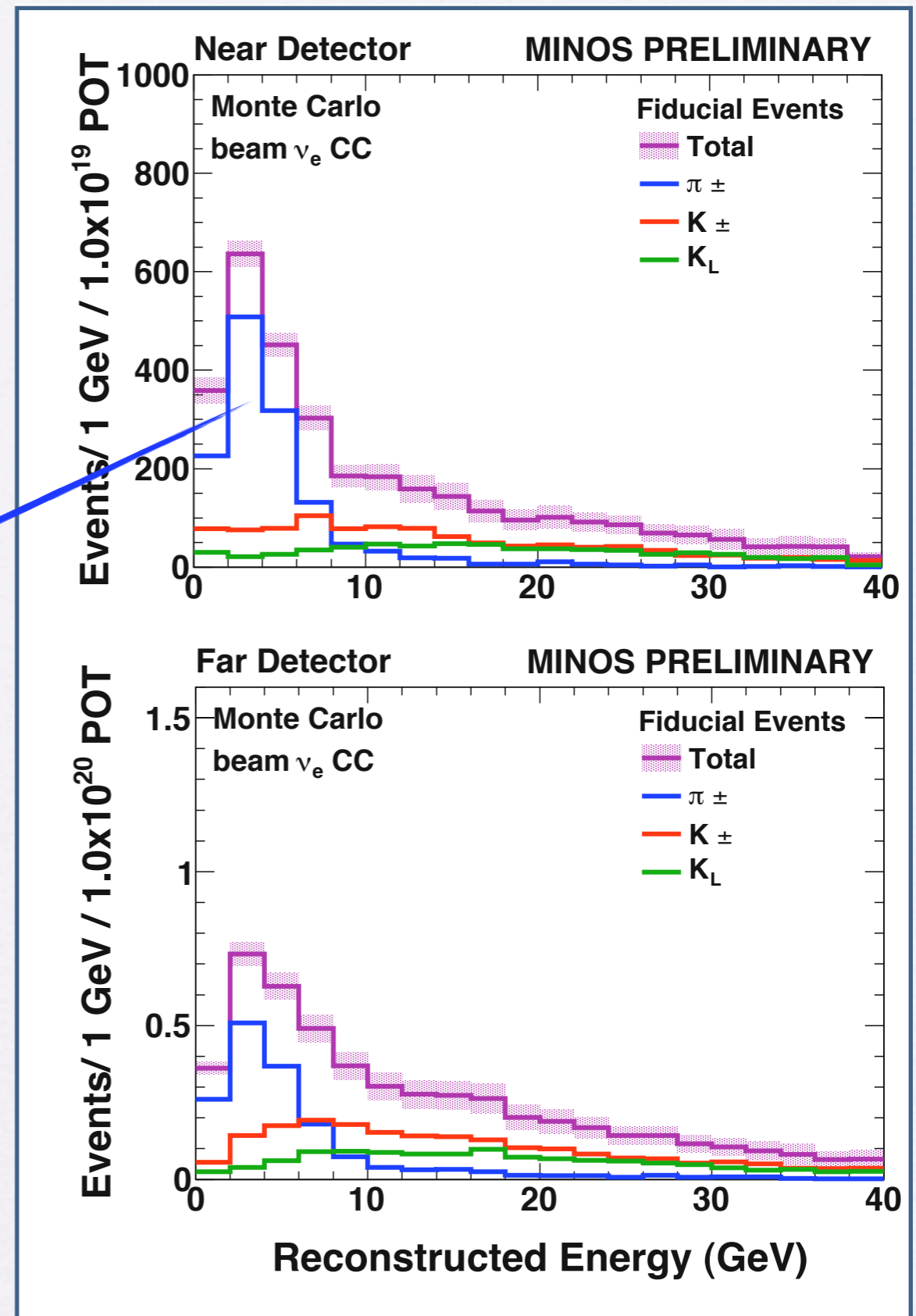
Backup

# Beam $\nu_e$ component

- Neutrino beam has 1.3% of  $\nu_e$  contamination from pion and kaon decays.
- Region of interest for the  $\nu_e$  oscillation analysis, 1-8GeV, dominated by events from secondary muon decays:

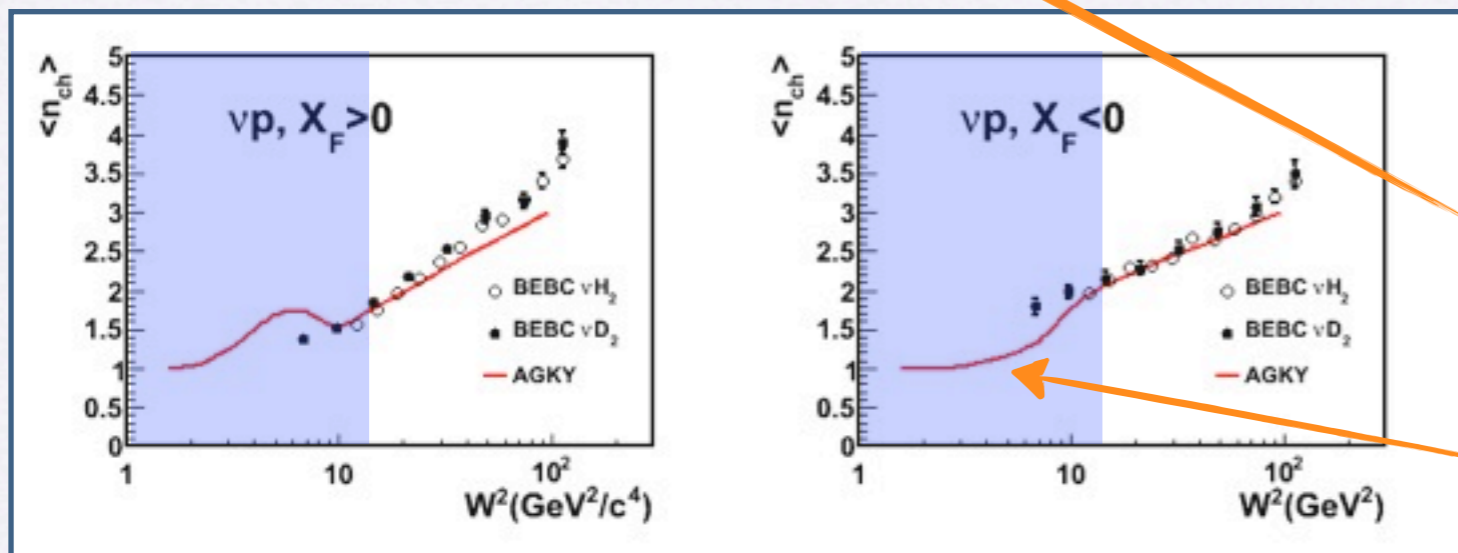
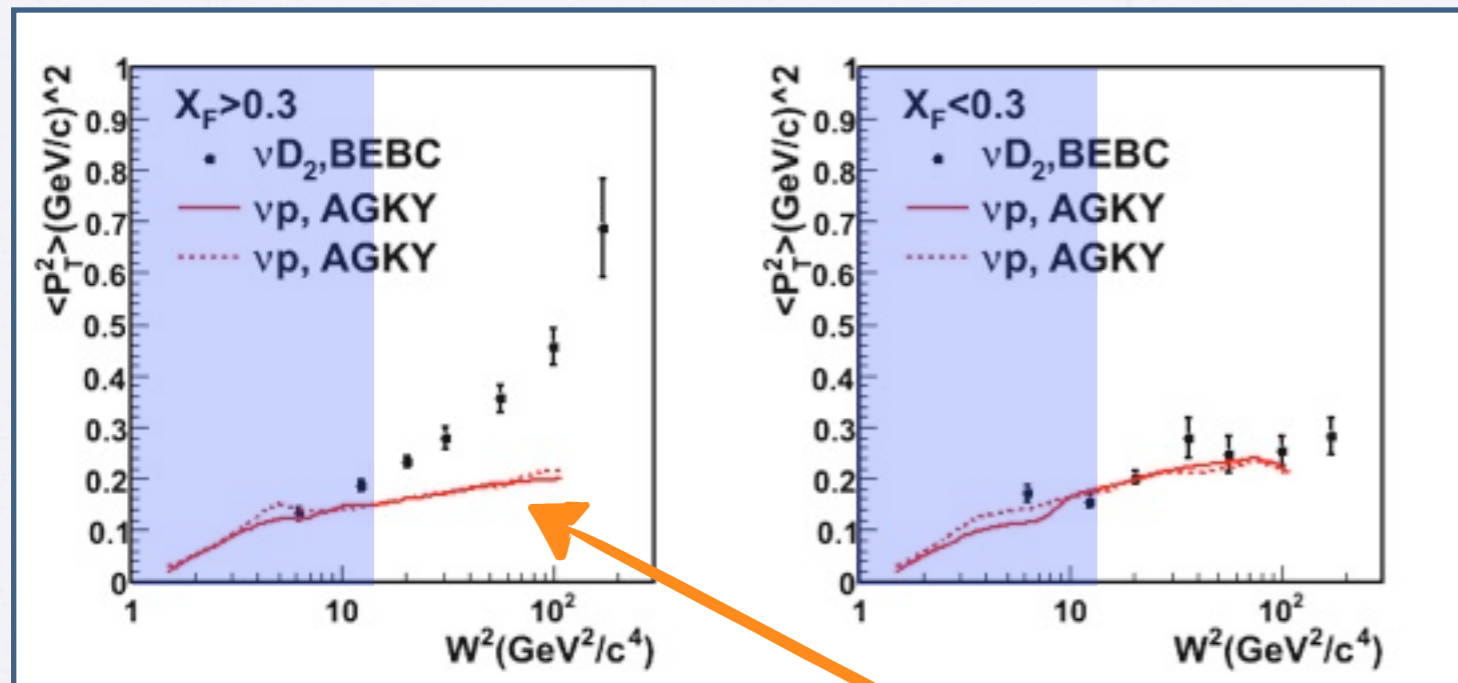


- Near and Far beam  $\nu_e$  spectra are constrained by using  $\nu_\mu$  events with different beam configurations.
- Errors from these fits after  $\nu_e$  selection are  $\sim 9\%$  in the Near and  $\sim 13\%$  in the Far Detector.





# MINOS Monte Carlo



Region of interest: 1 - 15  $\text{GeV}^2$  in  $W^2$

**We need to use more information from our own data in the Near Detector.**

- MC tuned to external bubble chamber data for hadronization models.
- Tuning focused in the following quantities:
  - Charged/neutral pion multiplicity and dispersion.
  - Forward/backward fragments.
  - Fragmentation functions.
  - Transverse momentum.
- Transverse momentum still too low in forward hemisphere.
- Model at lower  $W^2$  is an extrapolation.

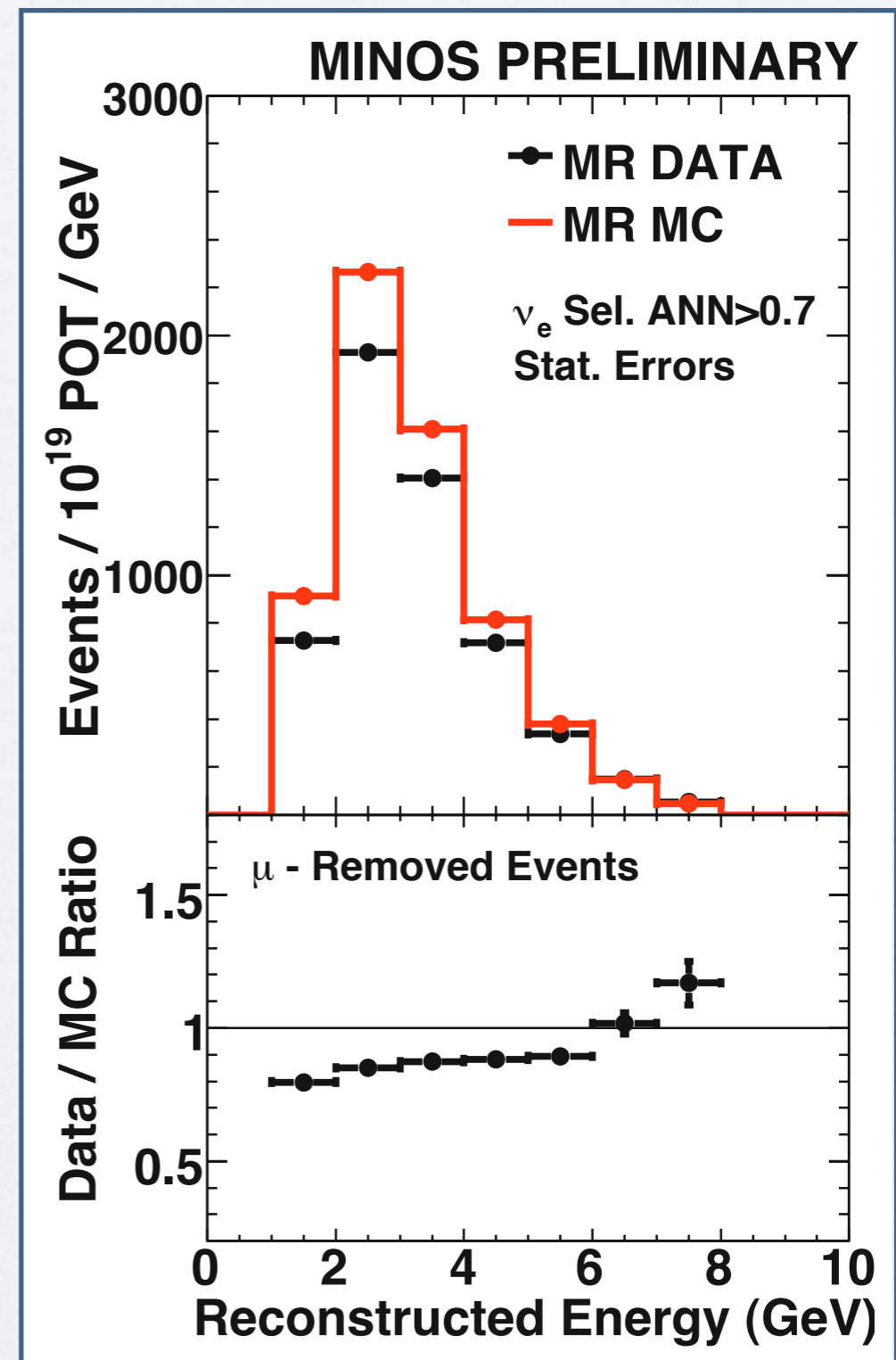
# Using MRCC as a data-driven correction

- We use the data/MC ratio from MRCC to obtain a **data-driven correction** that is applied to the standard NC events as a function of energy.

$$NC_i^{corr} = \frac{MRCC_i^{data}}{MRCC_i^{MC}} \times NC_i^{MC}$$

- The number of  $\nu_\mu$  CC events is taken from the number of events in the data minus the corrected NC and beam  $\nu_e$  events.
- Differences between NC and MRCC showers introduces a systematic error that is difficult to quantify.

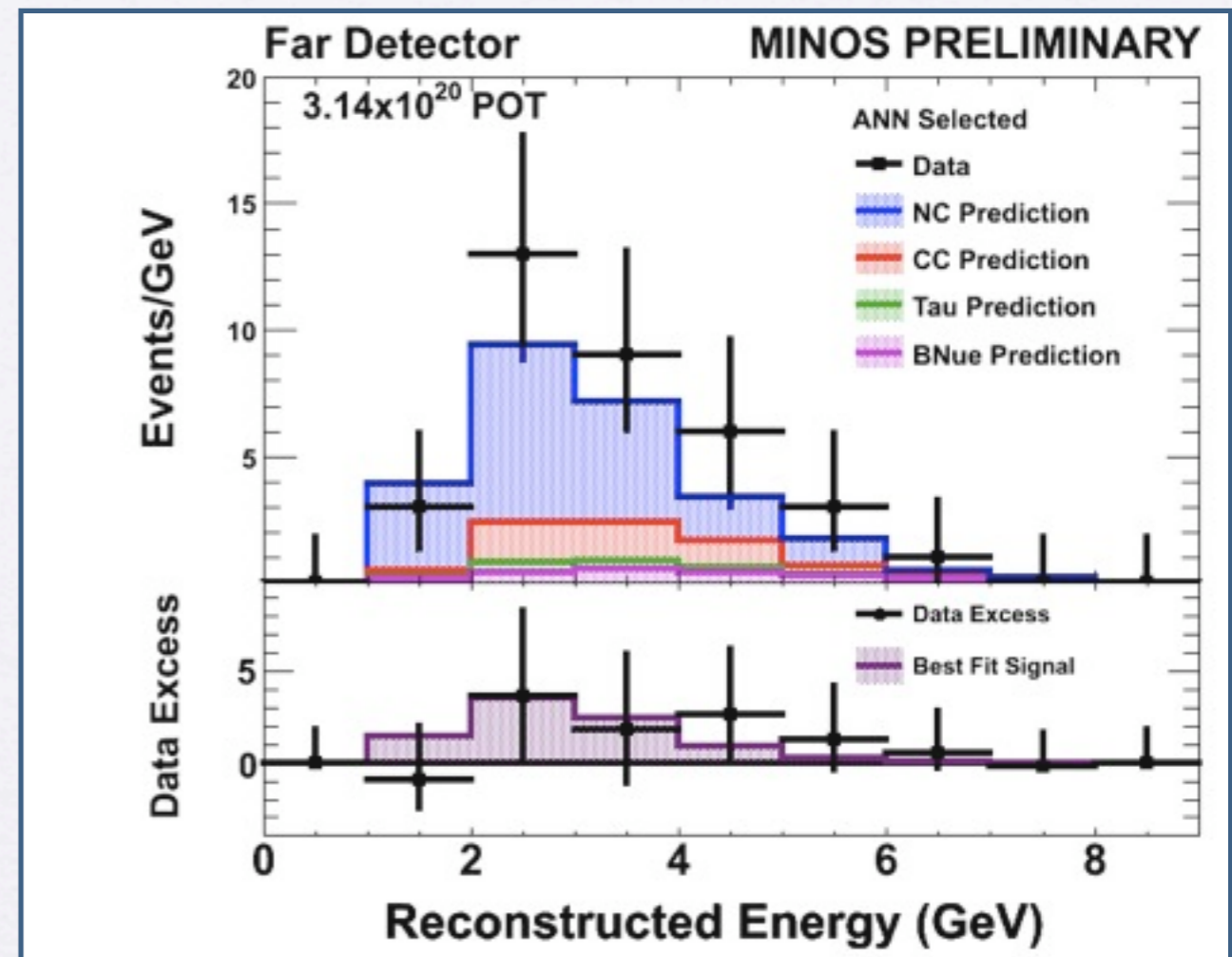
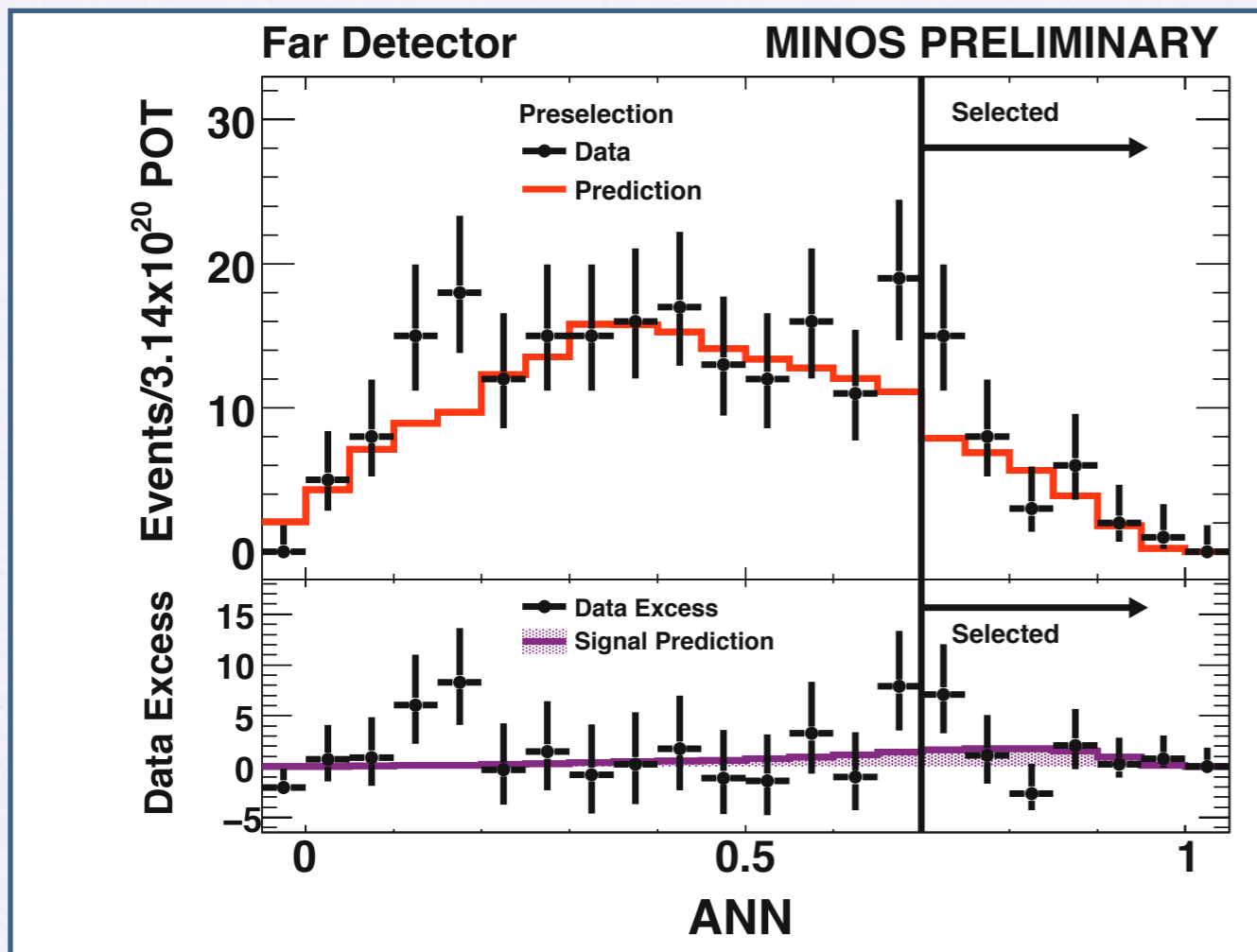
## Secondary separation method



# $\nu_e$ Selected Far Detector Data

## Primary selection method

- We observe a total of 35 events in this sample.
- We expect  $27 \pm 5(\text{stat}) \pm 2(\text{sys})$  background events.

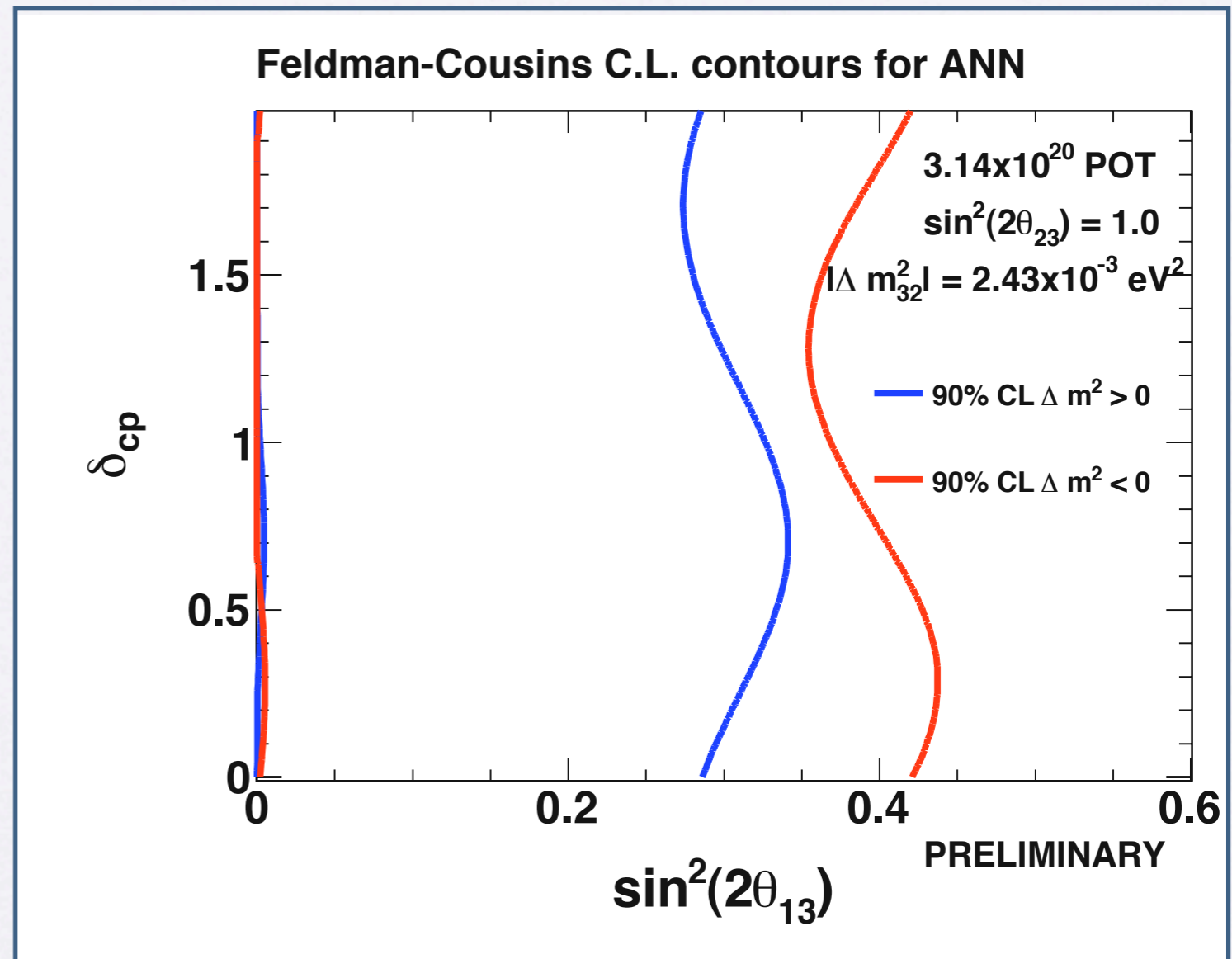


- If we fit the oscillation hypothesis to data, we can obtain the signal prediction for the best fit point.

# MINOS 90% CL in $\sin^2 2\theta_{13}$

Fitting the oscillation hypothesis to our data

- Plot shows 90% limits in  $\delta_{CP}$  vs.  $\sin^2 2\theta_{13}$ 
  - shown at the MINOS best fit value for  $\Delta m^2_{32}$  and  $\sin^2 2\theta_{23}$ .
  - for both mass hierarchies
- A Feldman-Cousins method was used.
- Results are for primary selection and primary separation method.



# Searching for $\theta_{13}$

Missing element in the PNMS neutrino mixing matrix

- The probability of  $\nu_e$  appearance in a  $\nu_\mu$  beam:

$$P(\nu_\mu \rightarrow \nu_e) \approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \Delta$$

$$\alpha \equiv \Delta m_{21}^2 / \Delta m_{31}^2$$

$$\Delta \equiv \frac{\Delta m_{31}^2 L}{4E}$$

$$+\alpha \Delta \cos \theta_{13} \sin 2\theta_{13} \cos \delta \sin 2\theta_{12} \sin 2\theta_{23} \sin \Delta \cos \Delta$$

$$-\alpha \Delta \cos \theta_{13} \sin 2\theta_{13} \sin \delta \sin 2\theta_{12} \sin 2\theta_{23} \sin \Delta \sin \Delta$$

no matter effects

- Searching for  $\nu_e$  events in MINOS, we can access  $\sin^2(2\theta_{13})$ .
- Probability depends not only on  $\theta_{13}$  but also on  $\delta_{CP}$ .
  - A non-zero  $\theta_{13}$  would open the door to a CP violation measurement in the neutrino sector which could reveal the origin of the matter/anti-matter asymmetry of the universe.

# Searching for $\theta_{13}$

Adding matter effects

- The probability of  $\nu_e$  appearance in a  $\nu_\mu$  beam:

$$A \equiv \frac{G_f n_e L}{\sqrt{2}\Delta} \approx \frac{E}{11 \text{ GeV}}$$

$$P(\nu_\mu \rightarrow \nu_e) \approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2(A-1)\Delta}{(A-1)^2} + 2\alpha \sin \theta_{13} \cos \delta \sin 2\theta_{12} \sin 2\theta_{23} \frac{\sin A\Delta}{A} \frac{\sin(A-1)\Delta}{(A-1)} \cos \Delta - 2\alpha \sin \theta_{13} \sin \delta \sin 2\theta_{12} \sin 2\theta_{23} \frac{\sin A\Delta}{A} \frac{\sin(A-1)\Delta}{(A-1)} \sin \Delta$$
$$\Delta \equiv \frac{\Delta m_{31}^2 L}{4E}$$

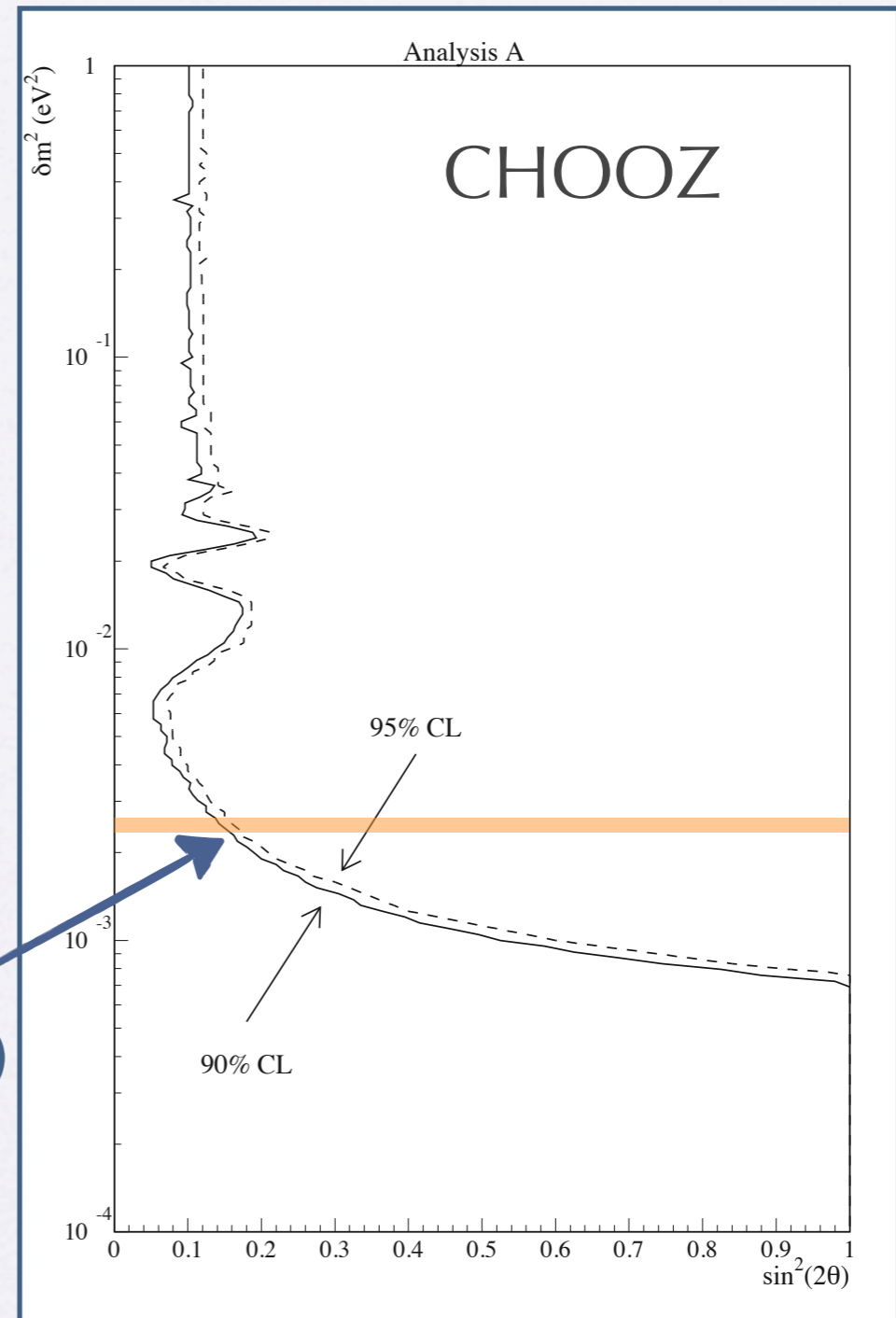
- Searching for  $\nu_e$  events in MINOS, we can access  $\sin^2(2\theta_{13})$ .
- Probability depends not only on  $\theta_{13}$  but also on  $\delta_{CP}$ .
- Probability is enhanced or suppressed due to **matter effects** which depend on the mass hierarchy i.e. the sign of  $\Delta m_{31}^2 \sim \Delta m_{32}^2$ .

# Relevant oscillation parameters

- The CHOOZ experiment published a limit in  $\sin^2(2\theta_{13})$ .
- Note reactor experiments do not have  $\delta_{CP}$  or mass hierarchy dependence.
- Since then MINOS has measured  $\Delta m^2_{32}$  very precisely.
- Thus for this talk:

**MINOS best fit**  
 $|\Delta m^2_{32}| = 2.43 \times 10^{-3} \text{ eV}^2$   
 $\sin^2 2\theta_{23} = 1.00$

**CHOOZ limit (90%CL)**  
 $\sin^2 2\theta_{13} = 0.15$



There are no measurements for  $\delta_{CP}$  or the mass hierarchy.

# Searching for $\theta_{13}$

Missing element in the PNMS neutrino mixing matrix

- The probability of  $\nu_e$  appearance in a  $\nu_\mu$  beam:

$$P(\nu_\mu \rightarrow \nu_e) \approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \sin^2 \Delta$$

$$\alpha \equiv \Delta m_{21}^2 / \Delta m_{31}^2$$

$$\Delta \equiv \frac{\Delta m_{31}^2 L}{4E}$$

$$+\alpha \Delta \cos \theta_{13} \sin 2\theta_{13} \cos \delta \sin 2\theta_{12} \sin 2\theta_{23} \sin \Delta \cos \Delta$$

$$-\alpha \Delta \cos \theta_{13} \sin 2\theta_{13} \sin \delta \sin 2\theta_{12} \sin 2\theta_{23} \sin \Delta \sin \Delta$$

no matter effects

- Searching for  $\nu_e$  events in MINOS, we can access  $\sin^2(2\theta_{13})$ .
- Probability depends not only on  $\theta_{13}$  but also on  $\delta_{CP}$ .
  - A non-zero  $\theta_{13}$  would open the door to a CP violation measurement in the neutrino sector which could reveal the origin of the matter/anti-matter asymmetry of the universe.



# Searching for $\theta_{13}$

Adding matter effects

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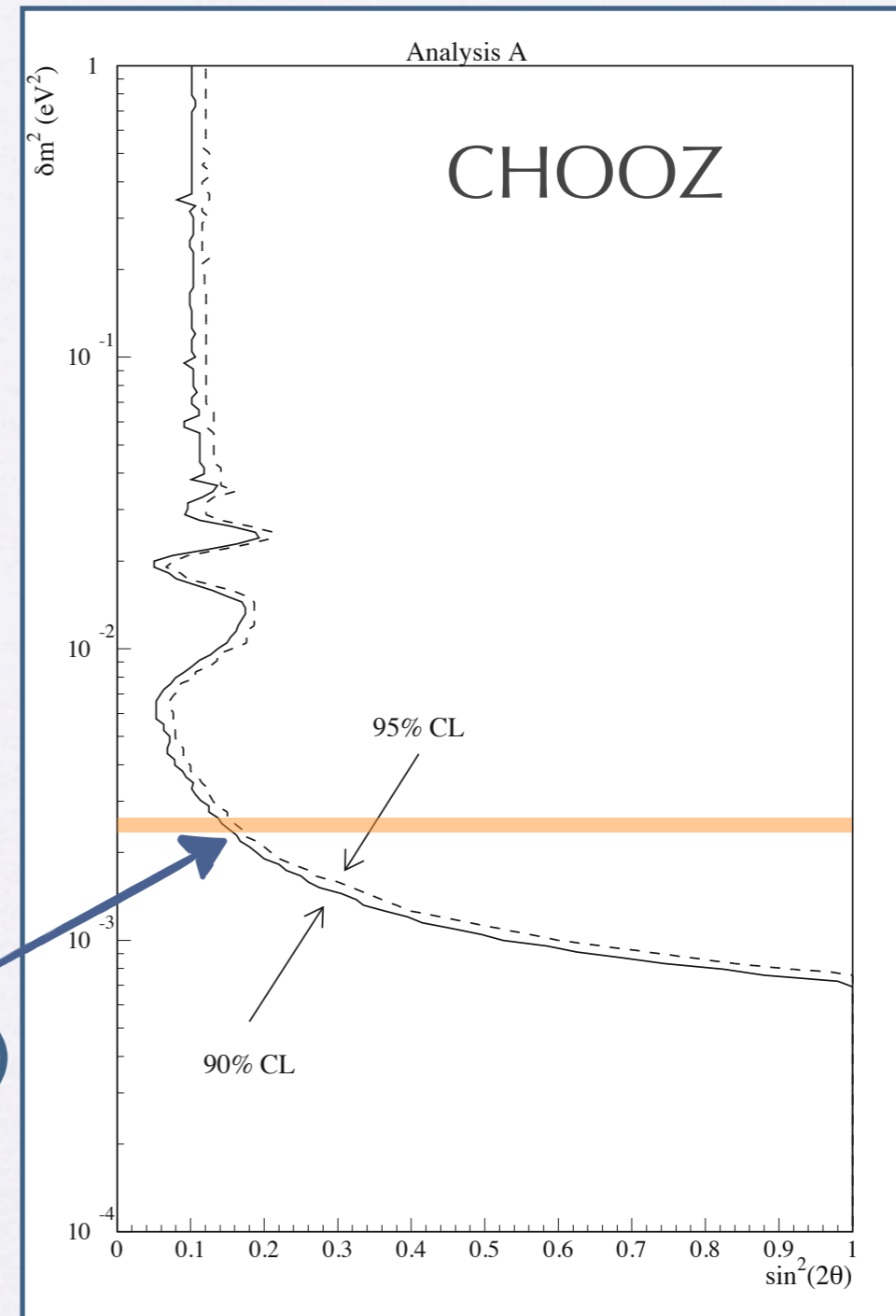
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